

# NON-SIMULTANEOUS QUENCHING IN A SYSTEM OF HEAT EQUATIONS COUPLED AT THE BOUNDARY

RAÚL FERREIRA

Dep. de Matemática Aplicada, U. Complutense de Madrid,

**Primer Encuentro de la Red de Ecuaciones  
Parabólicas y Elípticas No Lineales**

Joint work with

ARTURO DE PABLO, UC3M, Leganés, Spain.

MAYTE PÉREZ-LLANOS, UC3M, Leganés, Spain.

FERNANDO QUIRÓS, UAM, Madrid, Spain.

JULIO D. ROSSI, UBA, Buenos Aires, Argentina.

• **What is Quenching?** We say that the solution of a non-linear PDE quenches if

$$\|u_t(\cdot, T)\|_\infty = \infty \quad \text{while} \quad \|u(\cdot, T)\|_\infty < \infty.$$

Model problem [Kawarada'75]

$$\begin{cases} v_t = v_{xx} + \frac{1}{1-v}, & -L < x < L, 0 < t < T, \\ v(-L, t) = v(L, t) = 0, & 0 < t < T, \\ v(x, 0) = v_0(x), & -L < x < L. \end{cases}$$

Quenching happens when  $v$  reaches the level one and  $v_t \rightarrow \infty$ .

Moreover, if  $v_0$  is symmetric and decreasing for  $x > 0$ , then  $v$  reaches the level one only at  $x = 0$ . We denote that single-point quenching.

Other quenching problems:[Acker-Walter'78], [Levine'80–93], [Chan'96], [Galaktionov-Gerbi-Vázquez'99], [Fila-Guo'02].

## • Questions

- (1) What? The solution develop a singularity?
- (2) When? The singularity occurs at finite time  $T$  or not?
- (3) Where? We can determinate the point where the singularity happens?

$$Q(u) = \{x : \exists (x_n, t_n) \rightarrow (x, T) \text{ such that } u(x_n, t_n) \rightarrow 0\}.$$

is the **quenching set**.

- (4) How? Asymptotic behaviour.
  - **Quenching rate.**
  - **Asymptotic profile.**
- (5) After singularity. If we consider problem the problem as limit of approximated problems defined for every  $0 < t < \infty$ , we can study the possible continuation of the solution beyond  $t = T$ .
- (6) Numerical methods. We can find a numerical method which reproduces the same properties of the continuous solution?

## Problem

$$\left\{ \begin{array}{l} \left\{ \begin{array}{l} u_t = u_{xx}, \\ v_t = v_{xx}, \end{array} \right. \quad \text{in } (0, 1) \times (0, T), \\ \left\{ \begin{array}{l} u_x(0, t) = v^{-p}(0, t), \\ v_x(0, t) = u^{-q}(0, t), \end{array} \right. \quad u_x(1, t) = 0, \\ \left\{ \begin{array}{l} u(x, 0) = u_0(x), \\ v(x, 0) = v_0(x), \end{array} \right. \quad v_x(1, t) = 0, \end{array} \right. \quad \text{in } (0, T),$$

The initial data are  $C^2[0, 1]$  functions, increasing, concave and they satisfy the boundary condition.

## • A Escalar Quenching Problem

$$\begin{cases} u_t = u_{xx}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = u^{-p}(0, t), & 0 < t < T, \\ u_x(1, t) = 0, & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 < x < 1, \end{cases}$$

where  $p > 0$ .

- This problem has quenching in finite time,  $T$ , for all initial data.
- The quenching set is always de origin,  $Q(u) = \{0\}$ .
- The minimum of the solution verifies

$$u(0, t) \sim (T - t)^{\frac{1}{2(p+1)}}$$

- the behavior near the quenching time is given by a self-similar solution.
- For  $t > T$  we can define the continuation of the solution as a solution of

$$\begin{cases} \bar{u}_t = \bar{u}_{xx}, & 0 < x < 1, t > T, \\ \bar{u}(0, t) = 0, & t > T, \\ \bar{u}_x(1, t) = 0, & t > T, \\ \bar{u}(x, T) = u(x, T), & 0 < x < 1, \end{cases}$$

[Fila-Levine'93],[Christov-Deng'01],[Fila-Guo'02]

## What?

- Quenching happens in finite time for every initial data.

*Proof.*

-  $u \leq M = \|u_0\|_\infty$  and  $v \leq N = \|v_0\|_\infty$ .

- By Integration in space

$$\int_0^1 u(s, t) ds \leq M - V^{-p}(0)t,$$

$$\int_0^1 v(s, t) ds \leq N - U^{-q}(0)t.$$

□

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## When?

- $T \leq \min\{MV^p(0), NU^q(0)\}$

## What?

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□

## When?

- $T \leq \min\{MV^p(0), NU^q(0)\}$

## Where?

- The only quenching point is the origin.

*Proof.* Both variables are supersolution of the following problem

$$\begin{cases} w_t = w_{xx} & \text{in } (0, 1) \times (0, T) \\ w(0, t) = w(1, t) = 0 & \text{in } (0, T) \\ w(x, 0) = w_0(x) \leq \min\{u_0, v_0\} & \text{in } (0, 1) \end{cases}$$

□



## Preliminaries

- $u(\cdot, t)$  is an increasing function.
- $u(x, \cdot)$  is a decreasing function.
- There exists a positive constant such that

$$U'(t) \leq -CU^{-q}(t), \quad V'(t) \leq -CV^{-p}(t).$$

\* Therefore,

$$U(t) \geq C(T - t)^{1/(q+1)} \quad V(t) \geq C(T - t)^{1/(p+1)}$$

- There exists a positive constant such that

$$U'(t) \geq -CV^{-p-1}(t)U^{-q}(t), \quad V'(t) \geq -CU^{-q-1}(t)V^{-p}(t).$$

If  $u$  quenches while  $v$  remain positive then

$$U'(t) \sim -U^{-q}(t),$$

$$U(t) \sim (T - t)^{1/(q+1)}.$$

## Simultaneous vs Non-simultaneous

- If  $q < 1$  then for every  $v_0$  there exists  $u_0$  such that  $u$  quenches while  $v$  does not.

*Proof.* From the representation formula

$$\begin{aligned} \frac{1}{2}V(t) &= \int_0^1 \Gamma(y, t)v_0(y) dy + \int_0^t v(1, s) \frac{\partial \Gamma}{\partial x}(-1, t-s) ds \\ &\quad - \int_0^t U^{-q}(s)\Gamma(x, t-s) ds \end{aligned}$$

where  $\Gamma(x, t) = (4\pi t)^{-1/2}e^{-x^2/4t}$

- by apriori lower estimate for  $U$ , we obtain

$$V(t) \geq C_1 - C_2 \int_0^T (T-s)^{-\frac{q}{q+1}-\frac{1}{2}} = C_1 - C_2 T^{\frac{1-q}{2(1+q)}}$$

- From the upper estimate for  $T$ , we can choose  $U(0)$  small to conclude that  $V(t) \geq C_1/2$  for all  $0 \leq t \leq T$ .

□

## Simultaneous vs Non-simultaneous

- If  $q < 1$  then for every  $v_0$  there exists  $u_0$  such that  $u$  quenches while  $v$  does not.
- If  $v$  does not quench then  $q < 1$ .

*Proof.* From the representation formula

$$\begin{aligned} \frac{1}{2}V(t) &= \int_0^1 \Gamma(y, t)v_0(y) dy + \int_0^t v(1, s) \frac{\partial \Gamma}{\partial x}(-1, t - s) ds \\ &\quad - \int_0^t U^{-q}(s)\Gamma(x, t - s) ds \end{aligned}$$

- In the non-simultaneous case we have that

$$U(t) \sim (T - t)^{1/(1+q)}.$$

- Therefore,

$$V(t) \leq C_1 - C_2 \int_0^T (T - s)^{-\frac{q}{q+1} - \frac{1}{2}}.$$

But, the integral diverges if  $q \geq 1$ . □

## Simultaneous vs Non-simultaneous

- If  $q < 1$  then for every  $v_0$  there exists  $u_0$  such that  $u$  quenches while  $v$  does not.
- If  $v$  does not quench then  $q < 1$ .
- If  $p, q \geq 1$  then quenching is always simultaneous.
- If  $0 < p, q < 1$  there exist there exist initial data which produce simultaneous quenching.

*Proof.* Given  $(u_0, v_0)$ , we consider initial data  $(\lambda u_0, v_0)$ .

- From the representation formula and the lower estimates

$$V_\lambda(t) \geq C_1 - C_2 \int_0^{T_\lambda} (T_\lambda - s)^{-\frac{q}{q+1} - \frac{1}{2}} ds = C_1 - C_2 T_\lambda^{\frac{1-q}{2(1+q)}},$$

$$U_\lambda(t) \geq C_1 \lambda - C_2 \int_0^{T_\lambda} (T_\lambda - s)^{-\frac{p}{p+1} - \frac{1}{2}} ds = C_1 \lambda - C_2 T_\lambda^{\frac{1-p}{2(1+p)}},$$

- The estimate for the quenching time  $\Rightarrow T_\lambda \leq C \min\{\lambda^q, \lambda\}$

- For  $\lambda \ll 1$ ,

$$V_\lambda(t) \geq C_1 - C_2 \lambda^{\frac{1-q}{2(1+q)}} > 0$$

- For  $\lambda \gg 1$ ,

$$U_\lambda(t) \geq C_1 \lambda - C_2 \lambda^{\frac{q(1-p)}{2(1+p)}} > 0$$

□

## Simultaneous vs Non-simultaneous

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- If  $v$  does not quench then  $q < 1$ .
- If  $p, q \geq 1$  then quenching is always simultaneous.
- If  $0 < p, q < 1$  there exist there exist initial data which produce simultaneous quenching.
- If  $q < 1$  and  $p \geq p_0 = (1 + q)/(1 - q)$  then quenching is always non-simultaneous.

*Proof.* Assume that quenching is simultaneous.

- From the representation formula for  $v$ ,

$$0 = \int_0^1 \Gamma(y, T-t)v(y, t) dy + \int_t^T v(1, s) \frac{\partial \Gamma}{\partial x}(-1, T-s) ds - \int_t^T U^{-q}(s) \Gamma(0, T-s) ds$$

$$\text{Then, } V(t) \leq C \int_t^T U^{-q}(s)(T-s)^{-1/2} ds \leq C(T-t)^{(1-q)/2(1+q)}$$

- We introduce this upper estimate in the representation formula for  $u$

$$0 \leq \int_0^1 \Gamma(y, T)u(y, 0) dy + \int_0^T u(1, s) \frac{\partial \Gamma}{\partial x}(-1, T-s) ds - C \int_0^T (T-s)^{-\frac{1}{2} - \frac{p(1-q)}{2(1+q)}} ds$$

□

## Simultaneous vs Non-simultaneous

Summing up, we have that

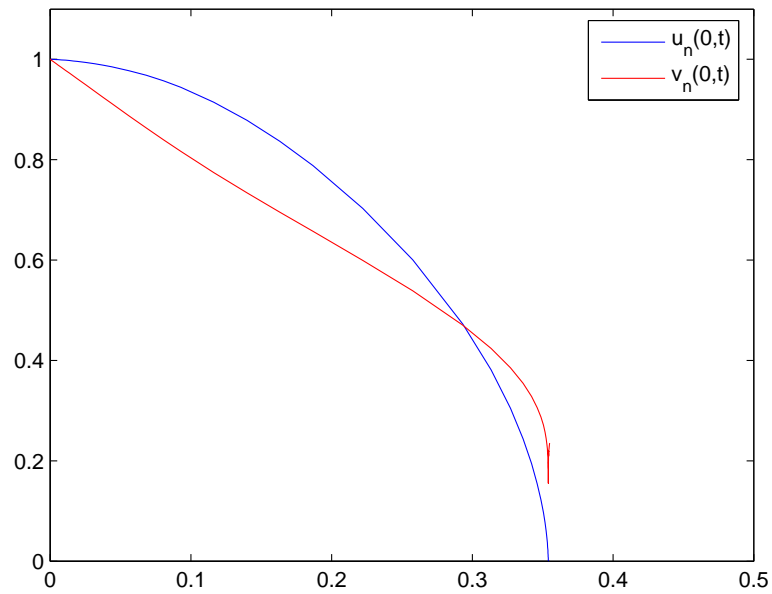
- If  $q < 1$  then for every  $v_0$  there exists  $u_0$  such that  $u$  quenches while  $v$  does not.
- If  $v$  does not quench then  $q < 1$ .
- If  $p, q \geq 1$  then quenching is always simultaneous.
- If  $0 < p, q < 1$  there exist there exist initial data which produce simultaneous quenching.
- If  $q < 1$  and  $p \geq p_0 = (1 + q)/(1 - q)$  then quenching is always non-simultaneous.

We conjecture that  $p_0 = 1$ .

## Simultaneous vs Non-simultaneous

- Numerical experiment with  $q = 1/3$  and  $p = 1$ . Initial data

$$u_0(x) = 1 + x, \quad v_0(x) = 1 + x - x^2.$$



## After Quenching.

We consider the problem  $(P_n)$

$$\left\{ \begin{array}{l} \left\{ \begin{array}{l} (u_n)_t = (u_n)_{xx}, \\ (v_n)_t = (v_n)_{xx}, \end{array} \right. \quad \text{in } (0, 1) \times (0, T), \\ \left\{ \begin{array}{l} (u_n)_x(0, t) = f_n(v(0, t)), \quad (u_n)_x(1, t) = 0, \\ (v_n)_x(0, t) = g_n(u(0, t)), \quad (v_n)_x(1, t) = 0, \end{array} \right. \quad \text{in } (0, T), \\ \left\{ \begin{array}{l} (u_n)(x, 0) = u_0(x), \\ (v_n)(x, 0) = v_0(x), \end{array} \right. \quad \text{in } (0, 1), \end{array} \right.$$

where

$$f_n(s) = \begin{cases} s^{-q}, & \text{if } s > 1/n, \\ n^{q+1}s, & \text{if } 0 < s \leq 1/n, \\ 0 & \text{if } s < 0, \end{cases}$$

$$g_n(s) = \begin{cases} s^{-p}, & \text{if } s > 1/n, \\ n^{p+1}s, & \text{if } 0 < s \leq 1/n, \\ 0 & \text{if } s < 0, \end{cases}$$

- The solution  $(u_n, v_n)$  are defined for all  $t > 0$ .
- A natural attempt to obtain a continuation of  $(u, v)$  after quenching is to pass to the limit as  $n \rightarrow \infty$  in  $(u_n, v_n)$ .
- This problem DOES NOT HAVE a Comparison Principle.



## After Quenching.

- $(u_n, v_n)$  is uniformly bounded from above.

*Proof.* Both components are a subsolution of the problem

$$\begin{cases} w_t = w_{xx}, & 0 < x < 1, t > 0, \\ w_x(0, t) = w_x(1, t) = 0, & t > 0, \\ w(x, 0) = \max(\|v_0\|_\infty, \|u_0\|_\infty), & 0 \leq x \leq 1. \end{cases}$$

□

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□

- $(u_n, v_n) > (u, v)$  for  $(x, t) \in [0, 1] \times [0, T)$ .

*Proof.*  $(u_n, v_n)$  is a supersolution to the original problem

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□

- Therefore, for  $t \in [0, T)$  there exists

$$(u_\infty, v_\infty) = \lim_{n \rightarrow \infty} (u_n, v_n).$$

- $(u_\infty, v_\infty) = (u, v)$  for all  $t < T$ .

## After non-simultaneous quenching.

- Let  $\tau_n$  the first time at truncation take place. Observe that  $\tau_n \rightarrow T$  as  $n \rightarrow \infty$ .

• [Key] For each  $n$  sufficiently large, there exists a time  $T_n$ , such that

$$K_1 \leq V_n(t) \leq K_2 \text{ for all } t > 0,$$

$$U_n(t) < 0 \text{ for all } t > T_n,$$

for some constants  $K_i > 0$  independent of  $n$ . Moreover, we have the estimate

$$\tau_n < T_n < \tau_n + C/n^2.$$

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for some constants  $K_i > 0$  independent of  $n$ . Moreover, we have the estimate

$$\tau_n < T_n < \tau_n + C/n^2.$$

•  $u_n$  is a supersolution to the problem

$$\begin{cases} w_t = w_{xx} & 0 < x < 1, t > 0, \\ w_x(0, t) = K_1^{-p}, & t > 0, \\ w_x(1, t) = 0, & t > 0, \\ w(x, 0) = u_0(x). \end{cases}$$

- Therefore  $(u_n, v_n)$  are uniformly bounded in compact sets and there exists

$$(u_\infty, v_\infty) = \lim_{n \rightarrow \infty} (u_n, v_n).$$

What problem verifies  $(u_\infty, v_\infty)$ ?

## After non-simultaneous quenching.

- Now we just observe that, for  $n$  large,  $(u_n, v_n)$  is a solution to

$$\begin{cases} (u_n)_t = (u_n)_{xx}, & (v_n)_t = (v_n)_{xx}, & 0 < x < 1, t > T_n, \\ (u_n)_x(0, t) = (v_n)^{-p}(0, t), & (v_n)_x(0, t) = 0, & t > T_n, \\ (u_n)_x(1, t) = 0, & (v_n)_x(1, t) = 0, & t > T_n, \end{cases}$$

- Since  $\tau_n < T_n < \tau_n + C/n^2$ , we have that  $T_n \rightarrow T$

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- Since  $\tau_n < T_n < \tau_n + C/n^2$ , we have that  $T_n \rightarrow T$

### Passing to the limit in the $v_n$ component

- The integral version of the problem for  $v_n$  is for  $t > T_n$ ,

$$-\int_{T_n}^t \int_0^1 v_n \varphi_t + \int_0^1 v_n(t) \varphi(t) - \int_0^1 v_n(T_n) \varphi(T_n) = \int_{T_n}^t \int_0^1 v_n \varphi_{xx} - \int_{T_n}^t v_n \varphi_x \Big|_0^1$$

- Since  $K_1 < v < K_2$ , we can pass to the limit in all the terms. The only tricky point is to show that  $v_n(x, T_n) \rightarrow v(x, T)$ .

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### Passing to the limit in the $v_n$ component

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- Since  $K_1 < v < K_2$ , we can pass to the limit in all the terms. The only tricky point is to show that  $v_n(x, T_n) \rightarrow v(x, T)$ .

- Therefore,  $v_\infty$  verifies the problem

$$\begin{cases} (v_\infty)_t = (v_\infty)_{xx} & \text{in } (0, 1) \times (T, \infty) \\ (v_\infty)_x(0, t) = (v_\infty)_x(1, t) = 0 & \text{in } (T, \infty) \\ v_\infty(x, T) = v(x, T) & \text{in } (0, 1) \end{cases}$$



## After non-simultaneous quenching.

- $v_n(x, T_n) \rightarrow v(x, T)$  as  $n \rightarrow \infty$

*Proof.* Let  $G$  the Green function of the Neumann problem. Then, for  $\tau_n \leq t \leq T_n$

$$v_n(x, t) = \int_0^1 G(x-y, t)v(y, \tau_n) dy + \int_{\tau_n}^t G(x, t-s)(v_n)_x(0, s) ds$$

- The first integral converges uniformly to  $v(x, T)$
- The second integral tends to zero

$$\int_{\tau_n}^t G(x, t-s)(v_n)_x(0, s) ds \leq Cn^{q+1}(t-\tau_n) \leq Cn^{q+1}n^{-2} = Cn^{q-1} \rightarrow 0$$

□

## After non-simultaneous quenching.

### Passing to the limit in the $u_n$ component

- In this case the integral version reads for  $t > 0$

$$\begin{aligned} - \int_0^t \int_0^1 u_n \varphi_t + \int_0^1 u_n(t) \varphi(t) - \int_0^1 u_0 \varphi(0) \\ = \int_0^t \int_0^1 u_n \varphi_{xx} - \int_0^t V_n^{-p} \varphi - \int_0^t u_n \varphi_x \Big|_0^1 \end{aligned}$$

- Since  $v_n \geq K_1$ , we can pass to the limit in all terms to obtain that  $u_\infty$  verifies the problem

$$\left\{ \begin{array}{ll} (u_\infty)_t = (u_\infty)_{xx} & \text{in } (0, 1) \times (0, \infty) \\ (u_\infty)_x(0, t) = v_\infty^{-p}(0, t) & \text{in } (0, \infty) \\ (u_\infty)_x(1, t) = 0 & \text{in } (0, \infty) \\ u_\infty(x, 0) = u_0(x) & \text{in } (0, 1) \end{array} \right.$$

## After non-simultaneous quenching.

- Proof of Lemma [Key]

- At time  $t = \tau_n$  the functions  $u_n$  and  $v_n$  are increasing and concave. Therefore,

$$c \leq v_n(x, \tau_n) \leq v_n(0, \tau_n) + n^q x \leq C + n^q x ,$$
$$\frac{1}{n} \leq u_n(x, \tau_n) \leq \frac{1}{n} + (v_n)^{-p}(0, \tau_n) x \leq \frac{1}{n} + Cx .$$

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$$\begin{aligned} c &\leq v_n(x, \tau_n) \leq v_n(0, \tau_n) + n^q x \leq C + n^q x, \\ \frac{1}{n} &\leq u_n(x, \tau_n) \leq \frac{1}{n} + (v_n)^{-p}(0, \tau_n) x \leq \frac{1}{n} + Cx. \end{aligned}$$

- We estimate the time  $\widehat{\tau}_n$  at which  $v_n$  reaches the level  $c/2$ .

Denote by  $s(x, t) = v_n(x, t + \tau_n)$ , we have that  $s$  is super-solution to the problem

$$\left\{ \begin{array}{ll} h_t = h_{xx}, & 0 < x < 1, 0 < t < \infty, \\ h_x(0, t) = n^q, & 0 \leq t < \infty, \\ h_x(1, t) = 0, & 0 \leq t < \infty, \\ h(x, 0) = c, & 0 \leq x \leq 1. \end{array} \right.$$

This problem vanish in finite time. Let  $\tau_0$  be the time such that  $h(0, \tau_0) = c/2$ .

## After non-simultaneous quenching.

Rescaling  $h$  we take off the dependence on  $n$  in the boundary condition. Let

$$\psi(y, \tau) = h(y/n^q, \tau/n^{2q}).$$

which satisfies the problem

$$\begin{cases} \psi_\tau = \psi_{yy}, & 0 < y < n^q, 0 < \tau < \infty, \\ \psi_y(0, \tau) = 1, & 0 \leq \tau < \infty, \\ \psi_y(n^q, \tau) = 0, & 0 \leq \tau < \infty, \\ \psi(y, 0) = c, & 0 \leq y \leq n^q. \end{cases}$$

Then, there exists a time  $\tau_1$  at which  $\psi(0, \tau_1) = c/2$ . We have also that  $\psi(0, \tau_1) = h(0, \tau_1/n^{2q})$ , thus,

$$\widehat{\tau}_n \geq \tau_n + \tau_0 = \tau_n + \tau_1/n^{2q}.$$

## After non-simultaneous quenching.

- We estimate the time  $T_n$  at which  $u_n$  reaches the level 0.

While  $f_n(V_n(t)) = V_n^{-p}(t)$ , the function

$$r(x, t) = u_n(x, t + \tau_n)$$

is a subsolution to the linear problem

$$\left\{ \begin{array}{ll} r_t = r_{xx}, & 0 < x < 1, 0 < t < \infty, \\ r_x(0, t) = K^{-p}, & 0 < t \leq \infty \\ r_x(1, t) = 0, & 0 < t \leq \infty, \\ r(x, 0) = u(x, \tau_n), & 0 \leq x \leq 1, \end{array} \right.$$

where  $K = \max V_n(t)$ .

## After non-simultaneous quenching.

$$\left\{ \begin{array}{ll} r_t = r_{xx}, & 0 < x < 1, 0 < t < \infty, \\ r_x(0, t) = K^{-p}, & 0 < t \leq \infty \\ r_x(1, t) = 0, & 0 < t \leq \infty, \\ r(x, 0) = u(x, \tau_n), & 0 \leq x \leq 1, \end{array} \right.$$

This problem vanish in finite time  $\tau_0$ . Moreover,

$$r(0, t) \leq 0, \quad \text{for all } t > \tau_0.$$

$\omega(y, \tau) = n r(y/n, \tau/n^2)$ , verifies

$$\left\{ \begin{array}{ll} \omega_\tau = \omega_{yy}, & 0 < y < n, 0 < \tau < \infty, \\ \omega_y(0, \tau) = K^{-p}, & 0 \leq \tau < \infty, \\ \omega_y(n, \tau) = 0, & 0 \leq \tau < \infty, \\ \omega(y, 0) = nu(y/n, \tau_n), & 0 \leq y \leq n. \end{array} \right.$$

There exist  $\tau_1$  such that

$$0 = \omega(0, \tau_1) = nr(0, \tau_1/n^2)$$

and

$$\omega(0, \tau) \leq 0, \quad \text{for } \tau \geq \tau_1.$$

## After non-simultaneous quenching.

- Summing up we have

i) The time  $\widehat{\tau}_n$  at which  $v_n$  reaches the level  $c/2$  verifies that

$$\widehat{\tau}_n \geq \tau_n + C/n^{2q}.$$

ii) the time  $T_n$  at which  $r$  reaches the level 0 verifies that

$$T_n \leq \tau_n + C/n^2$$

iii) Since  $q < 1$ , we have that  $T_n \leq \widehat{\tau}_n$ .

iv) Therefore at time  $T_n$ ,  $u_n$  vanishes while  $v_n$  remains positive.

v)  $v_n > c/2$  for times greater than  $T_n$ .

vi) As  $v_n > c/2$  for  $t > 0$ , we have that  $u_n < r$  for all time, and

$$u_n(0, t) \leq 0, \quad \text{for all } t > \tau_n + C/n^2.$$



## Numerical Approximations

- We consider uniform mesh and its associated standard approximation of the second derivative.

$$\begin{cases} u'_1 = \frac{2}{h^2}(u_2 - u_1) - \frac{2}{h}v_1^{-p}, \\ u'_k = \frac{1}{h^2}(u_{k-1} - 2u_k + u_{k+1}), \\ u'_{N+1} = \frac{2}{h^2}(u_N - u_{N+1}), \\ u_k(0) = u_0(x_k), \quad k = 1, \dots, N + 1, \end{cases}$$

$$\begin{cases} v'_1 = \frac{2}{h^2}(v_2 - v_1) - \frac{2}{h}u_1^{-q}, \\ v'_k = \frac{1}{h^2}(v_{k-1} - 2v_k + v_{k+1}), \\ v'_{N+1} = \frac{2}{h^2}(v_N - v_{N+1}), \\ v_k(0) = v_0(x_k), \quad k = 1, \dots, N + 1. \end{cases}$$

• This method converges in set of the form  $[0, 1] \times [0, T - \tau]$  for all  $\tau > 0$ ,

$$\max_{t \in [0, T - \tau]} \max_k \{|u(x_k, t) - u_k(t)|, |v(x_k, t) - v_k(t)|\} \leq Ch$$

• Both functions,  $u_h$  and  $v_h$ , are increasing in space and decreasing in time. In fact, there exists a positive constant

$$u'_1(t) < -C, \quad v'_1(t) < -C.$$

Then,

$$u_1(t) \geq (T_h - t), \quad v_1 \geq (T_h - t).$$

# Numerical Approximations

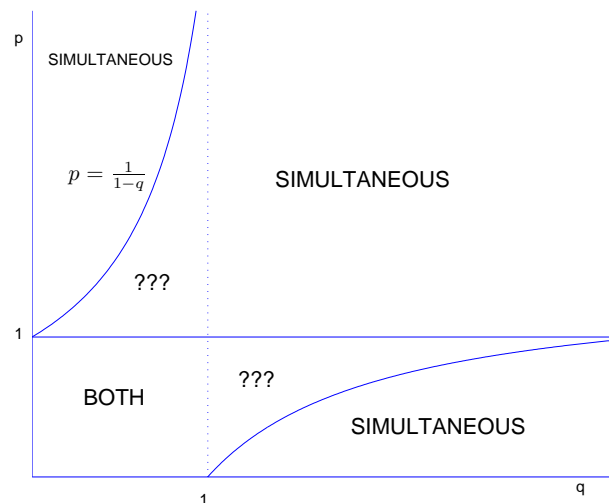
- This method quenches in finite time  $T_h$ . Moreover,

$$T_h \leq (1 + h) \min\{v_1(0), v_2(0)\}$$

and

$$T_h \rightarrow T.$$

- The only quenching point is the origin.
- Simultaneous vs Non-simultaneous



- In non-simultaneous case, the quenching rate is

$$u_1(t) \sim (T_h - t)$$

## Numerical Approximations

- We consider an adaptive method. We impose that

$$c_1 \leq -u_1^q u_1' \leq c_2$$

which is equivalent to

$$c_1 \leq R(t, h) := u_1^q \left( \frac{2}{h} v_1^{-p} - \frac{2}{h^2} (u_2 - u_1) \right) \leq c_2$$

- Let  $t_1$  be the time at which  $R$  reaches the tolerance  $c_1$ , at this point we refine the mesh.

We move the point  $x_2$  to a new place  $z$  while the rest of the mesh remain fixed, and we choose  $(u_z, v_z)$ , that is, the value of  $(u_h, v_h)$  at that new point  $z$ , such that

$$\frac{u_2(t_1) - u_1(t_1)}{h} = \frac{u_z(t_1) - u_1(t_1)}{z}, \quad \frac{v_2(t_1) - v_1(t_1)}{h} = \frac{v_z(t_1) - v_1(t_1)}{z},$$

i.e., the points  $(0, u_1(t_1))$ ,  $(z, u_z(t_1))$ ,  $(h, u_2(t_1))$  lay in the same line joining  $(0, u_1(t_1))$ ,  $(h, u_2(t_1))$ .

## Numerical Approximations

- At this time,

$$\begin{aligned} R(z, t_1) &= \left( \frac{2}{z} v_1^{-p}(t) - \frac{2}{z^2} (u_z(t) - u_1(t)) \right) u_1^q(t) \\ &= \frac{1}{z} \left( 2v_1^{-p}(t) - \frac{2}{z} (u_z(t) - u_1(t)) \right) u_1^q(t) \\ &= \frac{1}{z} \left( 2v_1^{-p}(t) - \frac{2}{h} (u_2(t) - u_1(t)) \right) u_1^q(t) \\ &= \frac{h}{z} R(h, t_1) > R(h, t_1) = c_1. \end{aligned}$$

So, we apply again the method with the new mesh up to time  $t_2$  at which  $R(z, t_2) = c_1$ . At this time we refine the mesh .....