

Stability of Unduloids under the Volume Preserving Mean Curvature Flow

David Hartley

Instituto de Ciencias Matemáticas

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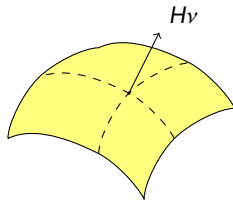
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$$\tilde{\mathbf{X}} : M^n \rightarrow \tilde{\mathbf{X}}(M^n) =: \tilde{\Omega} \subset \mathbb{R}^{n+1}$$

$$\frac{\partial \mathbf{X}}{\partial t} = (h - H) \nu, \quad \mathbf{X}(\cdot, 0) = \tilde{\mathbf{X}}, \quad h = \frac{\int_{M^n} H d\mu}{\int_{M^n} d\mu} \quad (1)$$



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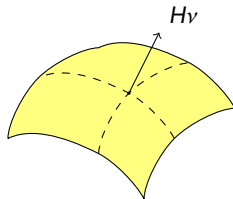
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- Together with boundary condition (if applicable)

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Graphs Over Hypersurfaces - $\frac{\partial \mathbf{x}}{\partial t} = (h - H) \nu$, $\mathbf{x}(\cdot, 0) = \tilde{\mathbf{x}}$

Considering the case where the initial hypersurface is a normal graph over an embedded hypersurface

$$\tilde{\mathbf{X}}(\mathbf{p}) = \mathbf{X}_{\rho_0}(\mathbf{p}) = \mathbf{X}_0(\mathbf{p}) + \rho_0(\mathbf{p}) \nu_0(\mathbf{p}), \quad \mathbf{p} \in M^n, \quad \rho_0 : M^n \rightarrow \mathbb{R}$$

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The flow is equivalent to an equation for the height function

$$\frac{\partial \rho}{\partial t} = L(\rho)(h(\rho) - H(\rho)), \quad \rho(\cdot, 0) = \rho_0 \quad (2)$$

$$L(\rho) = \sqrt{1 + \tilde{g}_p^{ij} \nabla_i \rho \nabla_j \rho}, \quad h(\rho) = \frac{\int_{M^n} H(\rho) d\mu_\rho}{\int_{M^n} d\mu_\rho}$$

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- \tilde{g}_ρ^{ij} is the inverse of $(\delta_i^k + \rho h_i^k) g_{kl} (\delta_j^l + \rho h_j^l)$.

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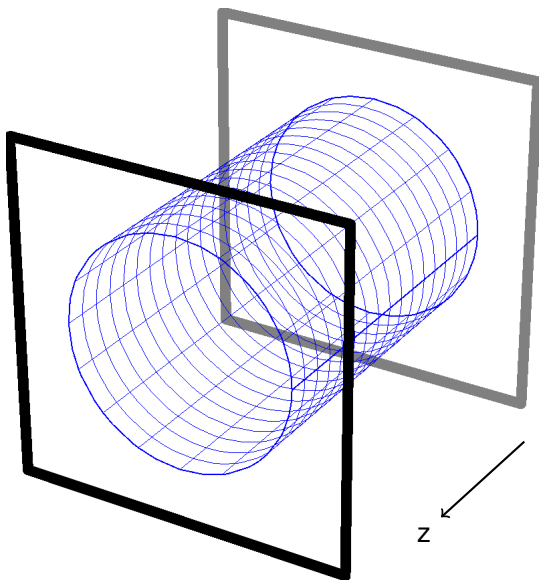
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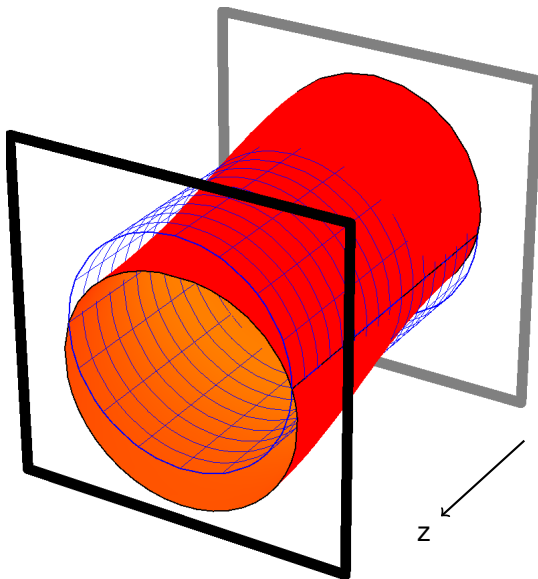
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Results for Cylinders - $\frac{\partial \rho}{\partial t} = L(\rho)(h(\rho) - H(\rho)), \rho(\cdot, 0) = \rho_0$

Theorem (H. 2014)

For all $\rho_0 \in U := \left\{ f \in h^{\frac{1,\beta}{\partial z}}(\mathcal{C}_{R,d}^n) : f > -R \right\}$ (2), together with the boundary condition $\frac{\partial \rho}{\partial z} \Big|_{z=0,d} = 0$, has a unique maximal solution

$$\rho \in C([0, \delta), U) \cap C^\infty((0, \delta), C^\infty(\mathcal{C}_{R,d}^n)).$$

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$$\rho \in C([0, \delta), U) \cap C^\infty((0, \delta), C^\infty(\mathcal{C}_{R,d}^n)).$$

Furthermore if $R > \frac{d\sqrt{n-1}}{\pi}$, then for each $k \in \mathbb{N}$ there exists a neighbourhood of zero $V_k \subset U$ such that for all $\rho_0 \in V_k$ we have $\delta = \infty$ and $\rho(t)$ converges exponentially fast, with respect to the C^k topology, to a function whose graph is a cylinder.

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- There exist non-axially symmetric hypersurfaces converging to cylinders

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Axial Symmetry- $\frac{\partial \rho}{\partial t} = L(\rho)(h(\rho) - H(\rho)), \rho(\cdot, 0) = \rho_0$

- In the axially symmetric case we have:

$$H(\rho) := \frac{-\frac{d^2 \rho}{dz^2}}{\left(1 + \left(\frac{d\rho}{dz}\right)^2\right)^{3/2}} + \frac{n-1}{\rho \sqrt{1 + \left(\frac{d\rho}{dz}\right)^2}}$$

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- To remove the boundary conditions we consider an equivalent equation on the circle of radius $\frac{d}{\pi}$:

$$\frac{\partial u}{\partial t} = G(u) := L(u)(h(u) - H(u)), \quad u(\cdot, 0) = u_0, \quad (3)$$

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$$\frac{\partial u}{\partial t} = G(u) := L(u)(h(u) - H(u)), \quad u(\cdot, 0) = u_0, \quad (3)$$

with solutions in the space $h_e^{2,\alpha} \left(\mathcal{S}_{\frac{d}{\pi}}^1 \right)$, where:

$$h_e^{2,\alpha} \left(\mathcal{S}_{\frac{d}{\pi}}^1 \right) = \left\{ u \in h^{2,\alpha} \left(\mathcal{S}_{\frac{d}{\pi}}^1 \right) : u(-z) = u(z) \right\}$$

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Splitting the Space - $\frac{\partial u}{\partial t} = L(u)(h(u) - H(u)), u(\cdot, 0) = u_0$

- We split the function space into $h_{e,0}^{2,\alpha}\left(\mathcal{S}_{\frac{d}{\pi}}^1\right) := P_0\left[h_e^{2,\alpha}\left(\mathcal{S}_{\frac{d}{\pi}}^1\right)\right]$,
 $P_0[u] := u - \int_{\mathcal{S}_{\frac{d}{\pi}}^1} u dz$, and the positive real line

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- The parameter is based on the volume due to its invariance under the flow

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- The parameter is based on the volume due to its invariance under the flow

There exist $U_1 \subset h_{e,0}^{2,\alpha}\left(\mathcal{S}_{\frac{d}{\pi}}^1\right) \times \mathbb{R}^+$ and $U_2 \subset h_e^{2,\alpha}\left(\mathcal{S}_{\frac{d}{\pi}}^1\right)$, neighbourhoods of $\left(0, \frac{\pi\sqrt{n-1}}{d}\right)$ and $\frac{d\sqrt{n-1}}{\pi}$, respectively, and a diffeomorphism $\psi : U_1 \rightarrow U_2$ such that:

$$\text{Vol}(\psi(\bar{u}, \eta)) = \frac{2d\omega_n(n-1)^n}{\eta^n}, \text{ and } \psi(\bar{u}, \eta) = \bar{u} + \int_{\mathcal{S}_{\frac{d}{\pi}}^1} \psi(\bar{u}, \eta) dz.$$

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$$\text{Vol}(\psi(\bar{u}, \eta)) = \frac{2d\omega_n(n-1)^n}{\eta^n}, \text{ and } \psi(\bar{u}, \eta) = \bar{u} + \int_{\mathcal{S}_{\frac{d}{\pi}}^1} \psi(\bar{u}, \eta) dz.$$

- The converse is also true: if $u \in U_2$ define $\bar{u} = P_0[u]$ and $\eta = (n-1)\left(\frac{2d\omega_n}{\text{Vol}(u)}\right)^{1/n}$ then $u = \psi(\bar{u}, \eta)$

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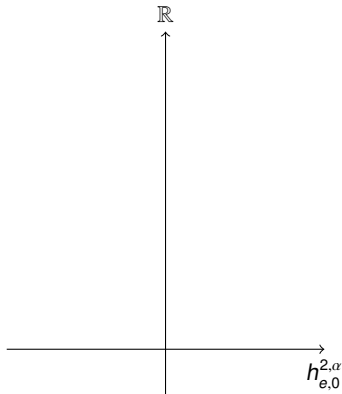
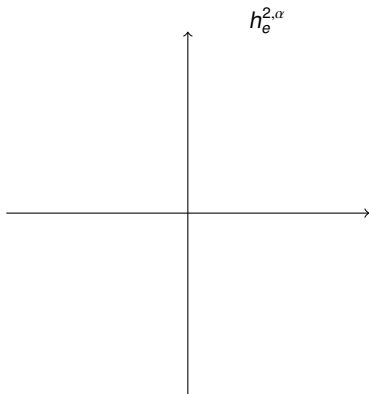
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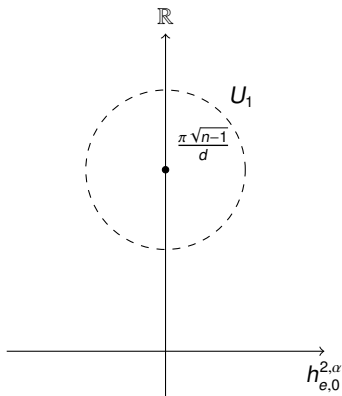
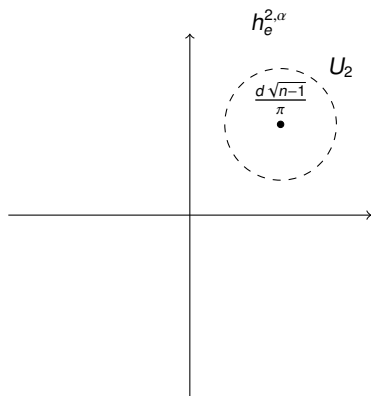
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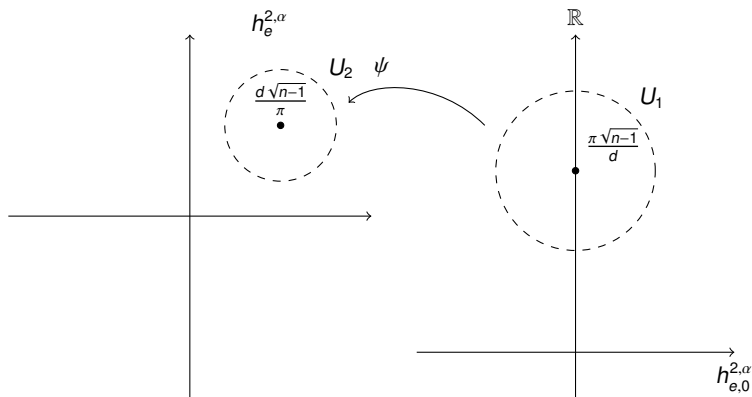
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Equivalence - $\frac{\partial u}{\partial t} = L(u)(h(u) - H(u)), u(\cdot, 0) = u_0$

$$\begin{aligned}\frac{\partial \bar{u}}{\partial t} &= \bar{G}(\bar{u}, \eta) := P_0 \left[G \left(\psi(\bar{u}(t), \eta) \right) \right], \\ \bar{u}(\cdot, 0) &= \bar{u}_0\end{aligned}\tag{4}$$

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- Solving this system is equivalent to solving (3) with $u_0 = \psi(\bar{u}_0, \eta)$, the solution is $u = \psi(\bar{u}, \eta)$

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- Solving this system is equivalent to solving (3) with $u_0 = \psi(\bar{u}_0, \eta)$, the solution is $u = \psi(\bar{u}, \eta)$
- A solution to (3) also gives a solution to (4), with η determined by the volume of u_0 , i.e. $\bar{u} = P_0[u]$

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Stationary Solutions - $\frac{\partial \bar{u}}{\partial t} = \bar{G}(\bar{u}, \eta), \bar{u}(\cdot, 0) = \bar{u}_0$

There exists $\delta > 0$ and a continuously differentiable curve of stationary solutions $(\bar{r}_s, \eta_s), |s| < \delta$, such that $(\bar{r}_0, \eta_0) = \left(0, \frac{\pi\sqrt{n-1}}{d}\right)$ and $\bar{r}_s \neq 0$ for $s > 0$. Further, there exists a neighbourhood of $\left(0, \frac{\pi\sqrt{n-1}}{d}\right)$ such that any non-trivial stationary solution in this neighbourhood is on the curve.

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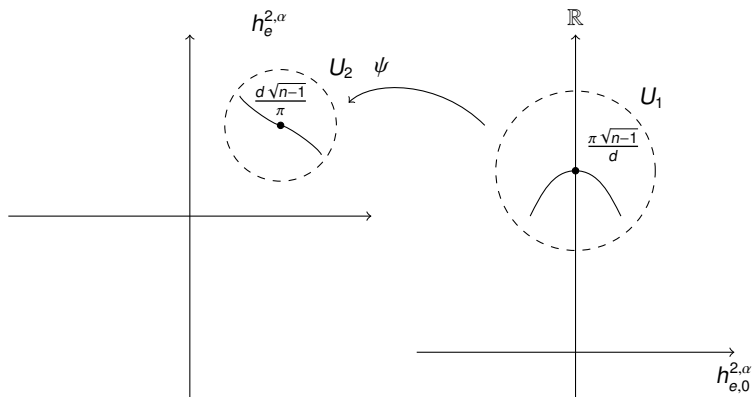
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Stationary Solutions - $\frac{\partial \bar{u}}{\partial t} = \bar{G}(\bar{u}, \eta), \bar{u}(\cdot, 0) = \bar{u}_0$



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- Further calculation gives $\frac{d\eta_s}{ds} \Big|_{s=0} = 0$,
 $\frac{d^2\eta_s}{ds^2} \Big|_{s=0} = C(n^2 - 10n - 2)$

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- For $n \leq 10$ subcritical bifurcation, for $n \geq 11$ supercritical bifurcation.
- If $n \geq 11$ there exists $\epsilon > 0$ such that $Re\left(\sigma\left(D_1 \bar{G}(\bar{r}_s, \eta_s)\right)\right) \subset \mathbb{R}^-$ for $0 < |s| < \epsilon$.

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For all $0 < |s| < \epsilon$ there exists $V_s \subset h_{e,0}^{2,\alpha}\left(\mathcal{S}^1_{\frac{d}{\pi}}\right)$, a neighbourhood of \bar{r}_s , and $\omega > 0$ such that if $\bar{u}_0 \in V_s$ then the solution to (4) with $\eta = \eta_s$ exists for all time and satisfies

$$\|\bar{u}(\cdot, t) - \bar{r}_s\|_{h^{2,\alpha}} \leq Ce^{-\omega t} \|\bar{u}_0 - \bar{r}_s\|_{h^{2,\alpha}}.$$

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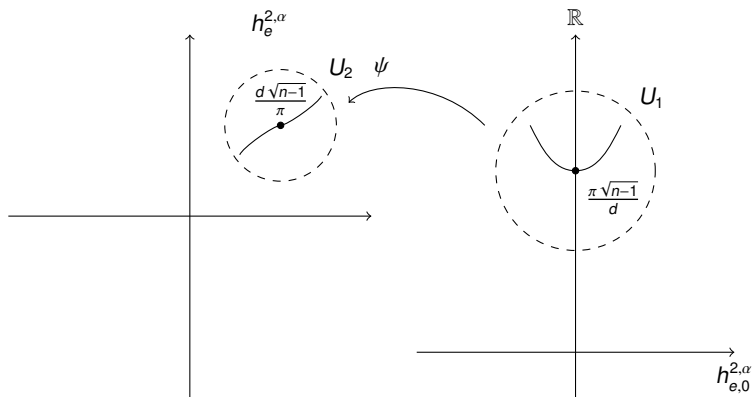
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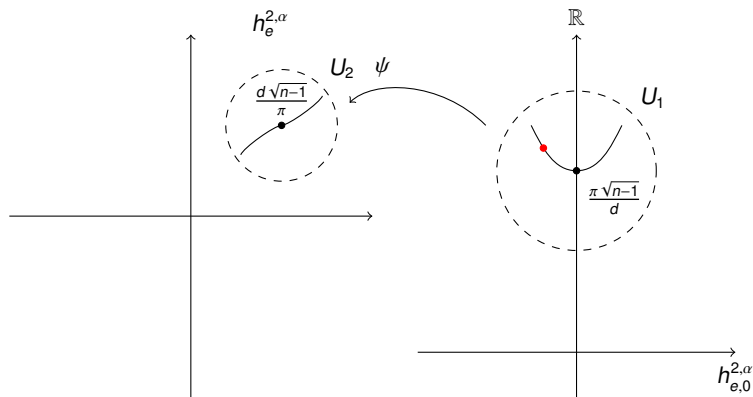
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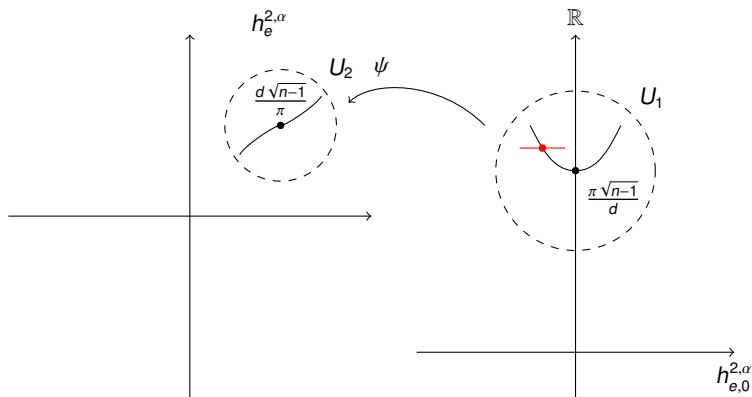
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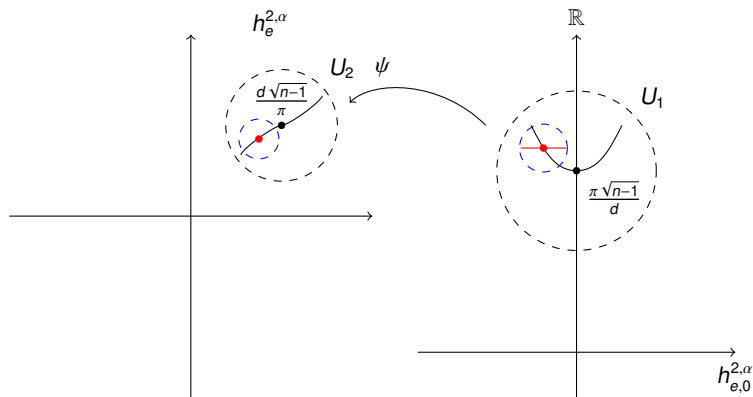
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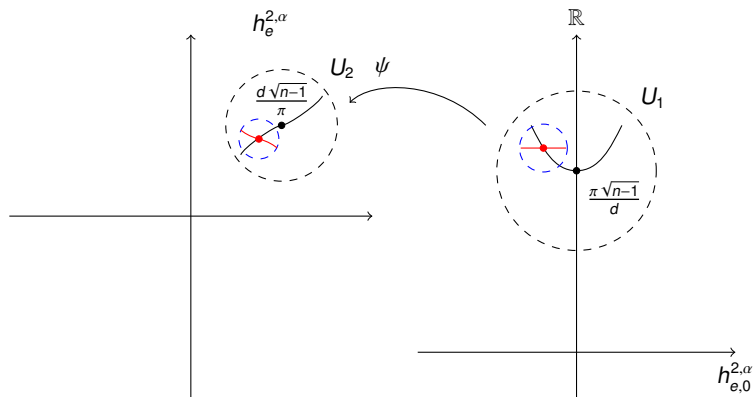
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Stable Unduloids - $\frac{\partial \rho}{\partial t} = L(\rho)(h(\rho) - H(\rho)), \rho(0) = \rho_0$

- $r_s := \psi(\bar{r}_s, \eta_s)|_{[0,d]}$, $|s| < \delta$, is a continuously differentiable family of profile curves that define CMC unduloids, with r_0 defining a cylinder of radius $\frac{d\sqrt{n-1}}{\pi}$.

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- For $n \leq 10$ unduloids close to the cylinder are unstable while for $n \geq 11$ unduloids close to the cylinder are stable under axially symmetric, volume preserving perturbations.

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- For $n \leq 10$ unduloids close to the cylinder are unstable while for $n \geq 11$ unduloids close to the cylinder are stable under axially symmetric, volume preserving perturbations.

Theorem (H. 2015)

When $n \geq 11$, for all $0 < |s| < \epsilon$ there exists $U_s \subset h^2_{d,\alpha}([0,d])$ a neighbourhood of r_s and $\omega > 0$ such that if $\rho_0 \in U_s$ and $\text{Vol}(\rho_0) = \text{Vol}(r_s)$ then the solution to (2) exists for all time and satisfies

$$\|\rho(\cdot, t) - r_s\|_{h^{2,\alpha}} \leq Ce^{-\omega t} \|P_0[u_{\rho_0} - u_{r_s}]\|_{h^{2,\alpha}}.$$

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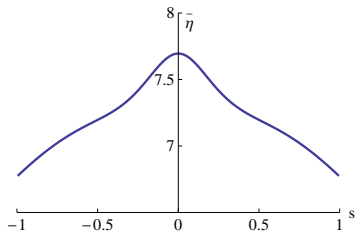
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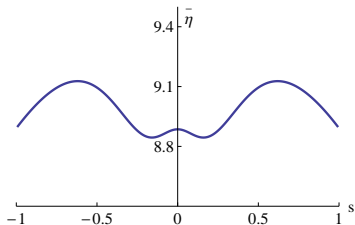
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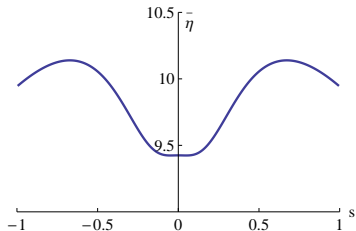
Bifurcation Parameter Curves



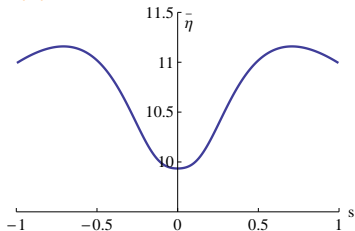
(a) Normalised η_s when $n = 7$



(b) Normalised η_s when $n = 9$



(c) Normalised η_s when $n = 10$



(d) Normalised η_s when $n = 11$

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Thank You

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