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# Information-theoretical complexity for the hydrogenic abstraction reaction 

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(Received 4 April 2011; final version received 26 May 2011)


#### Abstract

In this work, we have investigated the complexity of the hydrogenic abstraction reaction by means of information functionals such as disequilibrium $(D)$, exponential entropy $(L)$, Fisher information $(I)$, power entropy $(J)$ and joint information-theoretic measures, i.e. the $I-D, D-L$ and $I-J$ planes and the Fisher-Shannon and López-Mancini-Calbet (LMC) shape complexities. The analysis of the information-theoretical functionals of the one-particle density was computed in position $(r)$ and momentum $(p)$ space. The analysis revealed that all of the chemically significant regions can be identified from the information functionals and most of the informationtheoretical planes, i.e. the reactant/product regions $(\mathrm{R} / \mathrm{P})$, the transition state (TS), including those that are not present in the energy profile such as the bond cleavage energy region (BCER), and the bond breaking/forming regions $(B-B / F)$. The analysis of the complexities shows that, in position as well as in the joint space, the energy profile of the abstraction reaction bears the same information-theoretical features as the $L M C$ and $F S$ measures. We discuss why most of the chemical features of interest, namely the BCER and $B-B / F$, are lost in the energy profile and that they are only revealed when particular information-theoretical aspects of localizability $(L$ or $J)$, uniformity $(D)$ and disorder $(I)$ are considered.


Keywords: statistical complexity; Fisher information; information theory; chemical reactions

## 1. Introduction

In recent years there has been increasing interest in applying complexity concepts to study physical, chemical and biological phenomena. Complexity measures are understood as general indicators of pattern, structure, and correlation in systems or processes. Several alternative mathematical notions have been proposed for quantifying the concepts of complexity and information, including the Kolmogorov-Chaitin or algorithmic information theory [1], the classical information theory of Shannon and Weaver [2], Fisher information [3], and the logical [4] and the thermodynamical [5] depths, among others. Some of the above share rigorous connections with others as well as with Bayes and information theory [6]. The term complexity has been applied with different meanings: algorithmic, geometrical, computational, stochastic, effective, statistical, and structural, among others, and it has been employed in many fields: dynamical systems, disordered systems, spatial patterns, language, multielectronic systems, cellular automata, neuronal
networks, self-organization, DNA analyses, social sciences, astrophysics, among others [7,8].

The definition of complexity is not unique, and its quantitative characterization has been an important subject of research and it has received considerable attention $[9,10]$. The usefulness of each definition depends on the type of system or process under study, the level of the description, and the scale of the interactions among either elementary particles, atoms, molecules, biological systems, etc. Fundamental concepts such as uncertainty or randomness are frequently employed in the definitions of complexity, although other concepts such as clustering, order, localization or organization might also be important for characterizing the complexity of systems or processes. It is not clear how the aforementioned concepts might intervene in the definitions so as to quantitatively assess the complexity of the system. However, recent proposals have formulated this quantity as a product of two factors, taking into account order/ disequilibrium and delocalization/uncertainty. This is the case of the definition of the López-Mancini-Calbet

[^0](LMC) shape complexity [9-12] that, like others, satisfies the boundary conditions by reaching its minimal value in the extreme ordered and disordered limits. The $L M C$ complexity measure has been criticized [11], modified [12] and generalized [13], leading to a useful estimator that satisfies several desirable properties of invariance under scaling transfromations, translation, and replication [14,16]. The utility of this improved complexity has been verified in many fields [8] and allows reliable detection of periodic, quasiperiodic, linear stochastic, and chaotic dynamics [14-16]. The LMC measure is constructed as the product of two important information-theoretic quantities (see below): the so-called disequilibrium $D$ (also known as selfsimilarity [17] or information energy [18]), which quantifies the departure of the probability density from uniformity [12,15] (equiprobability), and the Shannon entropy $S$, which is a general measure of randomness/ uncertainty of the probability density [2], and quantifies the departure of the probability density from localizability. Both global quantities are closely related to the measure of spread of a probability distribution.

The Fisher-Shannon product FS has been employed as a measure of atomic correlation [19] and is also defined as a statistical complexity measure [20]. The product of the power entropy $J$ - explicitly defined in terms of the Shannon entropy (see below) - and the Fisher information measure $I$, combines both the global character (depending on the distribution as a whole) and the local character (in terms of the gradient of the distribution) to preserve the general complexity properties. Compared with the LMC complexity, apart from the explicit dependence on the Shannon entropy which serves to measure the uncertainty (localizability) of the distribution, the Fisher-Shannon complexity replaces the disequilibrium global factor $D$ by the Fisher local factor to quantify the departure of the probability density from disorder [3] of a given system through the gradient of the distribution.

The Fisher information $I$ itself plays a fundamental role in different physical problems, such as the derivation of the non-relativistic quantum-mechanical equations by means of the minimum $I$ principle [3], as well as the time-independent Kohn-Sham equations and the time-dependent Euler equation [21]. More recently, the Fisher information has also been employed as an intrinsic accuracy measure for specific atomic models and densities [22] as well as for general quantum-mechanical central potentials [23]. The concept of phase-space Fisher information has been analysed for hydrogen-like atoms and the isotropic harmonic oscillator [24]. Several applications concerning atomic distributions in position and momentum space have been performed where the $F S$ complexity
is shown to provide relevant information on atomic shell structure and ionization processes [20]. The Fisher measure has also been employed to test a density-based quantification of the steric effect of the ethane molecule [25].

On the other hand, theoretical chemistry has witnessed a great deal of research to study the energetics of chemical reactions [26]. For instance, a variety of calculations of potential energy surfaces have been performed at various levels of sophistication [27]. Within the broad scope of these investigations, particular interest has been focused on extracting information about the stationary points of the energy surface. Despite the fact that minima, maxima, and saddle points are useful mathematical features of the energy surface for following reaction paths [28], it has been difficult to attribute too much chemical or physical meaning to these critical points [29]. Whereas the reaction rate and the reaction barrier are chemical concepts that have been rigorously defined and experimentally studied since the early days of transition state (TS) theory [30], the structure of the TS remains a quest of physical organic chemistry. Understanding the TS is a fundamental goal of chemical reactivity theories, and implies knowledge of the chemical events that take place to better understand the kinetics and the dynamics of a reaction. On the other hand, a variety of density descriptors have been employed to study chemical reactions [30,31]. Among these, it is worth mentioning the reaction force studies on the potential energy of reactions that have been employed to characterize changes in the structural and/or electronic properties of chemical reactions [32,33]. In recent years, there has been increasing interest in the analysis of the electronic structure of atoms and molecules by applying information theory (IT) [34]. In recent studies we have shown that information-theoretic measures are capable of providing simple pictorial chemical descriptions of atoms and molecules. For instance, theoreticinformation analyses have shown to be useful for the phenomenological description of the course of elementary chemical reactions through the localized/ delocalized behavior of the electron densities in position and momentum space by revealing important chemical regions that are not present in the energy profile such as those in which bond forming and bond breaking occur and also the bond cleavage energy regions (BCER) [35]. Furthermore, the synchronous reaction mechanism of a $\mathrm{S}_{\mathrm{N}}$ 2-type chemical reaction and the non-synchronous behavior of the simplest hydrogen abstraction reaction were predicted by use of Shannon entropies analysis [36]. Also, the chemical phenomenon of $B-B / F$ was recently studied by the

Fisher information measure for both reactions, showing that this local measure in momentum space is highly sensitive in detecting these chemical events, whereas that in position space is able to detect differences in their mechanisms [37].

The goal of the present study is to perform an information-theoretical analysis of the hydrogenic abstraction reaction by use of information-theoretical measures and planes as well as the $L M C$ and $F S$ complexity products. Focus will be set on the recognition of patterns of uncertainty/localizability, disorder/ narrowness and disequilibrium/uniformity through $S, I$ and $D$, respectively. The organization of the paper is as follows. in Section 2 we defined the complexity measures along with their information-theoretic components. In Section 3 we calculate the information components as well as the Fisher-Shannon and LMC complexities. These information functionals of the oneparticle density are computed in position ( $r$ ), momentum $(p)$ as well as in the joint product space $(r p)$ that contains more complete information about the system. In addition, the Fisher-Shannon $(I-J)$ and the disequi-librium-Shannon $(D-L)$ and the Fisher-disequilibrium ( $I-D$ ) planes are studied. In Section 4, some conclusions are given.

## 2. Information-theoretical measures and complexities

In the independent-particle approximation, the total density distribution in a molecule is a sum of the contributions from the electrons in each of the occupied orbitals. This is the case in both $r$-space and $p$-space, position and momentum, respectively. In momentum space, the total electron density, $\gamma(\vec{p})$, is obtained through the molecular momentals (momentum-space orbitals) $\varphi_{i}(\vec{p})$, and similarly for the position-space density, $\rho(\vec{r})$, through the molecular position-space orbitals $\phi_{i}(\vec{r})$. The momentals can be obtained by three-dimensional Fourier transformation of the corresponding orbitals (and conversely)

$$
\begin{equation*}
\varphi_{i}(\mathbf{p})=(2 \pi)^{-3 / 2} \int \mathrm{~d} \mathbf{r} \exp (-i \mathbf{p} \cdot \mathbf{r}) \phi_{i}(\mathbf{r}) \tag{1}
\end{equation*}
$$

Standard procedures for the Fourier transformation of position space orbitals generated by ab-initio methods have been described [38]. The orbitals employed in $a b-$ initio methods are linear combinations of atomic basis functions and since analytic expressions are known for the Fourier transforms of such basis functions [39], the transformation of the total molecular electronic wavefunction from position to momentum space is computationally straightforward [40].

As mentioned in the introduction, the $L M C$ complexity is defined through the product of two relevant information-theoretic measures. For a given probability density in position space, $\rho(\vec{r})$, the $C(L M C)$ complexity is given by [9-12]

$$
\begin{equation*}
C_{r}(L M C)=D_{r} \mathrm{e}^{S_{r}}=D_{r} L_{r} \tag{2}
\end{equation*}
$$

where $D_{r}$ is the disequilibrium [17,18],

$$
\begin{equation*}
D_{r}=\int \rho^{2}(\mathbf{r}) \mathrm{d} \mathbf{r} \tag{3}
\end{equation*}
$$

and $S$ is the Shannon entropy [2],

$$
\begin{equation*}
S_{r}=-\int \rho(\mathbf{r}) \ln \rho(\mathbf{r}) \mathrm{d}^{3} \mathbf{r} \tag{4}
\end{equation*}
$$

from which the exponential entropy $L_{r}=\mathrm{e}^{S_{r}}$ is defined. Similar expressions for the LMC complexity measure in the conjugated momentum space might be defined for a distribution $\gamma(\vec{p})$

$$
\begin{equation*}
C_{p}(L M C)=D_{p} \mathrm{e}^{S p}=D_{p} L_{p} \tag{5}
\end{equation*}
$$

It is important to mention that the $L M C$ complexity of a system must comply with the following lower bound [41]:

$$
\begin{equation*}
C(L M C) \geq 1 \tag{6}
\end{equation*}
$$

The $F S$ complexity in position space, $C_{r}(F S)$, is defined in terms of the product of the Fisher information [3],

$$
\begin{equation*}
I_{r}=\int \rho(\mathbf{r})|\vec{\nabla} \ln \rho(\mathbf{r})|^{2} \mathrm{~d}^{3} \mathbf{r} \tag{7}
\end{equation*}
$$

and the power entropy [20] in position space $J_{r}$,

$$
\begin{equation*}
J_{r}=\frac{1}{2 \pi \mathrm{e}} \mathrm{e}^{(2 / 3) S_{r}} \tag{8}
\end{equation*}
$$

which depends on the Shannon entropy defined above. Therefore, the FS complexity in position space is given by

$$
\begin{equation*}
C_{r}(F S)=I_{r} \cdot J_{r} \tag{9}
\end{equation*}
$$

and similarly

$$
\begin{equation*}
C_{p}(F S)=I_{p} \cdot J_{p} \tag{10}
\end{equation*}
$$

in momentum space.
Let us remark that the factors in the power Shannon entropy $J$ are chosen to preserve the invariance under scaling transformations, as well as the rigorous relationship [42]

$$
\begin{equation*}
C(F S) \geq n \tag{11}
\end{equation*}
$$

with $n$ being the space dimensionality, thus providing a universal lower bound to $F S$ complexity. The definition in Equation (8) corresponds to the particular case $n=3$, the exponent containing a factor $2 / n$ for arbitrary dimensionality.

It is worth noting that the aforementioned inequalities remain valid for distributions normalized to unity, which is the choice that it is employed throughout this work for the three-dimensional molecular case.

Apart from the analysis of the position and momentum information measures, we have considered it useful to study these magnitudes in the product $r p$ space, characterized by the probability density $f(\vec{r}, \vec{p})=\rho(\vec{r}) \gamma(\vec{p})$, where the complexity measures are defined as

$$
\begin{equation*}
C_{r p}(L M C)=D_{r p} L_{r p}=C_{r}(L M C) C_{p}(L M C) \tag{12}
\end{equation*}
$$

and

$$
\begin{equation*}
C_{r p}(F S)=I_{r p} J_{r p}=C_{r}(F S) C_{p}(F S) \tag{13}
\end{equation*}
$$

From the above two equations, it is clear that the features and patterns of both the $L M C$ and $F S$ complexity measures in the product space will be determined by those of each conjugated space. However, the numerical analyses carried out in the next section reveal that the momentum space contribution plays a more relevant role than the position space contribution.

## 3. Complexity analysis of the hydrogenic abstraction reaction

The electronic structure calculations performed in the present study were carried out with the Gaussian 03 suite of programs [43]. Reported TS geometrical parameters for the abstraction reaction were employed [44]. Calculations for the IRC were performed at the MP2 (UMP2 for the abstraction reaction) level of theory with at least 35 points for each of the directions (forward/reverse) of the IRC. A high level of theory and a well-balanced basis set (diffuse and polarized orbitals) were then chosen to determine all of the properties for the chemical structures corresponding to the IRC. Thus, the QCISD(T) method was employed in addition to the $6-311++\mathrm{G}^{* *}$ basis set, unless otherwise stated. The molecular information measures were $S, D, I$ and $J$, the information planes $(D-L),(I-J)$ and $(I-D)$ and the complexity measures were $C(L M C)$ and $C(F S)$. All information-theoretical quantities are calculated in position and momentum space for the IRC path of the abtsraction reaction and obtained by employing software developed in our laboratory along
with 3D numerical integration routines [45], and the DGRID suite of programs [40].

The reaction $\mathrm{H}^{\bullet}+\mathrm{H}_{2} \rightarrow \mathrm{H}_{2}+\mathrm{H}^{\bullet}$ is the simplest radical abstraction reaction involving a free radical (atomic hydrogen) as a reactive intermediate. This kind of reaction involves at least two steps $\left(\mathrm{S}_{\mathrm{N}} 1\right.$ reaction type): in the first step, a new radical (atomic hydrogen in this case) is created by homolysis, and in the second step the new radical recombines with another radical species. Such homolytic bond cleavage occurs when the bond involved is not polar and there is no electrophile or nucleophile at hand to promote heterolytic patterns. The bond-breaking process requires energy that should be dissipated by relaxing the structure at the TS. Evidence has been presented [36] that shows that the two-step mechanism observed for this type of reaction is completely characterized by the Shannon entropies in conjugated space through concerted but yet asynchronous behavior.

Our calculations for this reaction were performed at two different levels: the IRC was obtained at the UMP2/6-311G level, and all properties at the IRC were obtained at the $\operatorname{QCISD}(\mathrm{T}) / 6-311++\mathrm{G}^{* *}$ level of theory. As a result of the IRC, 72 points evenly distributed between the forward and reverse directions of the IRC were obtained. A relative tolerance of $1.0 \times 10^{-5}$ was set for the numerical integrations [45].

Insight into the structural features of the distributions in both spaces concerning the global spreading (delocalization) of the densities can be obtained through the Shannon entropies in conjugated space. The behavior of the densities related to their local changes is appropriately described by a measure such as the Fisher information [3]. This measure quantifies the pointwise concentration of the electronic probability cloud, for each space, by means of the gradient content of the electron distribution, thus revealing the changes of the density and providing a quantitative estimation of its oscillatory character (smoothness). It is worth mentioning that, due to the global or local character of the particular information measure to be employed for the analysis, it is clear that each is capable of a partial description of all chemical phenomena, i.e. detection of the $\mathrm{R} / \mathrm{P}$ (reactant/product complexes) regions, $\mathrm{B}-\mathrm{B} / \mathrm{F}$ (bond-breaking/forming) regions, BCER, TS, mechanistic behavior, etc. Therefore, in order to allow for a full characterization of a chemical process it would be necessary to perform a complexity analysis because these informa-tion-theoretical products provide complementary information sources, i.e. $D$ (departure from uniformity) with $L$ (departure from localizability) through the $C(L M C)$ measure and also $I$ (departure from disorder) with $J$ (departure from localizability) through
the $C(F S)$ measure (Equations (2), (5), (9) and (10)). In the following sections we employ global and local quantities to provide a complete description of the hydrogenic abstraction reaction.

### 3.1. Information measures

Figure 1 depicts the values for the Shannon entropies in conjugated space so as to provide a description of the abstraction reaction in terms of the localizability features of the densities in conjugated space which will be of utility for analysing its complexity behavior. Both $C(L M C)$ and $C(F S)$ employ this global measure from the definition of $L$ and $J$, respectively. Summarizing the analysis performed in Ref. [35], the phenomenological description of the reaction shows the following: as the molecular complex approaches the TS its position space density becomes localized (minima of the position space density) in preparation for bond rupture, a process that requires energy. This is revealed by the local maxima of the corresponding momentum densities, i.e. its delocalization indicates a local increase of the necesary kinetic energy for bond cleavage at the BCER (bond cleavage energy regions) as observed in Figure 1. The homolysis then provokes energy/density relaxation of the molecules towards the TS, which is also observed from the figure.

We have found it useful to describe the chemical process in terms of the local/global features of the distributions at the IRC through the local measure $I$ (departure from disorder) and through the global measure $D$ (departure from uniformity). Therefore, in Figures 2 and 3 we depict both the Fisher information and the desequilibrium measures, in position and momentum space, respectively. Alluding to a previous study [37], a phenomenological description of the chemical course shows the following local features. Fisher information in position space (Figure 2) shows maxima at the $\mathrm{R} / \mathrm{P}$ region, whereas it has a global minimum at the TS, i.e. as the reactant complex approaches the TS the disorder of the distributions in position space increases. This may be interpreted chemically as follows: as the reaction evolves, the structural changes in the distributions diminish, which means that the $\mathrm{R} / \mathrm{P}$ has larger changes than the TS , which is characterized by a symmetric distribution (the smoothest of all at the IRC) produced by the spin coupling of the hydrogenic radical species. It is interesting to note that the gradient of the distributions increases towards the BCER at -1.0 a.u., indicating the beginning of the bond-cleavage region. Continuing with the analysis from Figure 2, the global features of the chemical course of the reaction indicate that the


Figure 1. Shannon entropies in conjugated spaces $S_{r}$ (red triangles) and $S_{p}$ (blue circles) for the IRC of the abstraction reaction.


Figure 2. Fisher information measure (red triangles) and disequilibrium (blue circles) in position space for the IRC of the abstraction reaction.


Figure 3. Fisher information measure (red triangles) and disequilibrium (blue circles) in momentum space for the IRC of the abstraction reaction.
disequilibrium for the $\mathrm{R} / \mathrm{P}$ show lower values than the BCER and that the TS has a global minimum, i.e. the position space distributions are the least uniform at the BCER, whereas the largest uniformity is observed at the TS. Chemically, the reaction proceeds by deforming the densities in the $\mathrm{R} / \mathrm{P}$ region so as to reach maxima at the BCER, i.e. position space densities at the onset of the bond-cleavage regions are the least uniform. These regions are also associated with the most delocalized densities in momentum space
.

The chemical significance of these features is that the structure of the $\mathrm{R} / \mathrm{P}$ is deformed so as to reach a higher degree of uniformity at the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions, which are associated with highly localized densities in position space according to the Shannon entropies depicted in Figure 1. The process then reverts so as to increase the disequilibrium (less uniformity) to reach the TS. It is interesting to observe that both processes (bond breaking and TS forming in the forward direction of the IRC) are located within the energy-releasing region described above for the Fisher measure, which means that both processes require energy. Interestingly, the structural energetic change to reach the TS is more pronounced than that required for bond breaking. This is because this reaction is driven by a spin coupling mechanism as a unique source of change. In the second step of the reaction the process is inverted and the energy accumulation region shown by Fisher in Figure 3 is now employed to leave the structural equiprobability (uniformity) in position space ( $D$ in Figure 2) implied by the TS. This is shown in Figure 3 by passing from structural disequilibrium to uniformity in momentum space (maximum and minimum values for $D$ at the TS and bondforming regions, respectively). The process then reverts by releasing the energy necessary for the structural deformation of the distribution in momentum space by diminishing the uniformity. It is worth mentioning that the reaction (in momentum space) might be characterized in terms of uniformity through the disequilibrium measure in that its minima describe the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions and the local maximum is associated with the TS.

### 3.2. Information planes

In the search of any joint features of disorderuniformity $(I-D)$, uniformity-localizability $(D-L)$, and disorder-localizability $(I-J)$ for the chemical course of the reaction we have found it useful to plot the contribution of $I$ and $D$ to the $I-D$ plane, of $D$ and $L$ to the total $L M C$ complexity, and similarly of $I$ and $J$ to the $F S$ complexity.

Notwithstanding that not all information products are good candidates to form complexity measures, i.e. to preserve the desirable properties of invariance under scaling, translation and replication, we have found it interesting to study the plane $I-D$, with the purpose of analysing patterns of disorder-uniformity. In Figures 4 and 5 we give a phenomenological description of the reaction through the $I_{r}-D_{r}$ and $I_{p}-D_{p}$ planes, respectively. In position space (Figure 4) we note that, at the $\mathrm{R} / \mathrm{P}$, the order is maximum and, as the reaction develops, both the disorder and


Figure 4. Fisher-disequilibrium plane in position space for the IRC of the abstraction reaction.


Figure 5. Fisher-disequilibrium plane in momentum space for the IRC of the abstraction reaction.
the disequilibrium increase until the BCER is reached. At this point the disequilibrium is maximum and then both the disorder and the uniformity increase until the TS is reached. At the TS, both the order and the disequilibrium are minimum. It is worth noting that the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region is not present in this plane. A general observation from Figure 4 is that the order/uniformity ratio possesses a negative slope that is nearly linear from $\mathrm{R} / \mathrm{P}$ to BCER , so that both quantities decrease until the BCER is reached and, at the same time, the order/disequilibrium ratio has a positive slope, indicating that both quantities decrease so as to reach the TS. Chemically, one may observe from Figure 4 that the structural changes in position space diminish along the IRC as the reaction evolves and the reactant complex reaches the TS, judging by
the Fisher measure, then the process reverts and the structural changes increase so as to reach the product complex. This effect is accompanied by a distortion of the position space densities so as to decrease/increase (depending on the direction of the reaction) their uniformity with the bending located at the BCER as the reaction evolves from the reactant complex up to the TS and then the process reverts at this point, i.e. the distortions acquire opposite directions by augmenting/ reducing (depending on the direction of the reaction) its disequilibrium so as to reach the product complex with the same inflexion point at the BCER.

In momentum space the $\mathrm{R} / \mathrm{P}$ regions exhibit maximum values for $I$ and $D$ (order and disequilibrium) and, as the reaction develops (forward direction), both quantities diminish so at to reach the BCER, which posseses maximum disorder, then the order increases by reducing $D$ up to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region, which has a minimum $D$ value, towards the TS , which has maximum local order. The general observation from Figure 5 is that the disorder/uniformity ratio possesses a positive slope, which is remarkably linear from $\mathrm{R} / \mathrm{P}$ to BCER, so that both quantities increase until the BCER is reached and then the slope becomes larger at the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region with both quantities decreasing so as to reach the TS. In chemical terms, the structural changes that the momentum space distributions exert decrease as the reaction evolves (forward direction), which conveys an energy accumulation (see analysis of Figure 3) up to where the BCER is reached, and then the process reverts by releasing energy with an increase in structural changes up to the TS. These energetic effects are accompanied by structural distortions in momentum space so as to increase the uniformity when the energy is accumulated from R to BCER, and then, as the energy is released, the distortions exert two opposite changes from BCER to TS, as explained below. According to the Fisher measure, from the BCER to the TS, the reaction requires energy, and judging by the $D$ measure, the reaction involves two stages: bond breaking and spin coupling, as we have explained in connection with Figure 3. This is clearly observed when structural uniformity increases from BCER to a global minimum at $\mathrm{B}-\mathrm{B} / \mathrm{F}$, and from this point the uniformity decreases so as to reach the TS.

It is worth comparing Figures 4 and 5 with respect to the behavior of the complementarity of the different spaces in the course of the chemical reaction. In position space (Figure 4) we observe that the BCER is characterized by minimum uniformity, whereas in momentum space (Figure 5) this region characterizes a state of maximum disorder. In both cases, one can note that the BCER reflects an inflection point, although the behavior in position space is that


Figure 6. Disequilibrium-Shannon plane ( $D-L$ ) in position space (red triangles) for the IRC of the abstraction reaction on a double logarithmic scale. The lower bound is depicted by the blue line (Equation (6)).


Figure 7. Disequilibrium-Shannon plane ( $D-L$ ) in momentum space (red triangles) for the IRC of the abstraction reaction on a double logarithmic scale. The lower bound is depicted by the blue line (Equation (6)).
of reaching a state of maximum disorder and uniformity together, whereas in momentum space the course of the reaction points to a state of local minimum disorder at the TS.

Figures 6 and 7 show plots (on a doublelogarithmic scale) of the $D_{r}-L_{r}$ and $D_{p}-L_{p}$ planes for the chemical reaction. At this point it is worth mentioning that there is a rigorous lower bound to the associated $C(L M C)$ complexity, given by Equation (6), which is $C(L M C)=D \cdot L \geq 1$ for both spaces. From both figures we can see that the $D-L$ plane is clearly separated into two regions, according to the $D \cdot L \geq 1$ inequality (valid for position, momentum as well as product spaces), and the region below the line (equality) corresponds to the forbidden region. Parallel lines to this bound represent isocomplexity regions,
showing that an increase (decrease) in uncertainty, $L$, along them is compensated by a proportional decrease (increase) in disequilibrium, and greater deviations from this frontier are associated with greater $L M C$ complexities.

From Figure 6 we note that, from the $\mathrm{R} / \mathrm{P}$ to the BCER regions, the behavior is isocomplex with the $C(L M C)$ bound, whereas for the rest of the IRC it behaves in a more complex manner. The general observations are that the $\mathrm{R} / \mathrm{P}$ show maximum uncertainty (highly delocalized structures) and, as the reaction evolves, both uniformity and uncertainity decrease up to the BCER, which exhibits maximum disequilibrium. Then uncertainty follows its decreasing path, whereas uniformity increases up to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region, which exhibits minimum uncertainty (highly localized structures) up to the TS by reducing the disequilibrium at the expense of increasing uncertainty so as to reach a structure with maximum uniformity. The chemical analysis proceeds by noting that, as the reaction develops (forward direction), the position space structures become distorted by losing uniformity and gaining localizability up to the BCER in preparation for bond cleavage. Then, from the BCER up to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region, we also observe two stages of the mechanism: first, the structures acquire both greater localizability and greater uniformity for bond rupture at the $B-B / F$. In the second stage, from $B-B / F$ to the TS, spin coupling is achieved by gaining uniformity at the expense of augmenting uncertainty.

In momentum space (Figure 7) we observe fairly isocomplex behavior from $\mathrm{R} / \mathrm{P}$ to BCER . In addition, the $\mathrm{R} / \mathrm{P}$ are characterized by possessing maximum disequilibrium and, as the reaction evolves, uniformity increases by augmenting the uncertainty up to the BCER, which has structures with maximum delocalizability. Then the structures gain more uniformity at the expense of lowering its uncertainty at the $\mathrm{B}-\mathrm{B} / \mathrm{F}$, which has the maximum uniformity. From this region up to the TS, uncertainty keeps diminishing at the expense of losing uniformity. The energetic analysis of Figure 7 completes the chemical picture by noting that, from $\mathrm{R} / \mathrm{P}$ to BCER, the structural changes exert an increase in delocalizability, and hence energy is being accumulated. The opposite is observed from BCER up to TS, where the structural changes are those that augment its localizability, hence releasing energy. As mentioned above, two stages are also observed from BCER to $\mathrm{B}-\mathrm{B} / \mathrm{F}$ and from this region to the TS, i.e. for bond breaking, uncertainty decreases at the expense of gaining uniformity, and for spin coupling in the second stage, it not only requires a decrease of uncertainty but also a reduction in uniformity.

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Figure 8. Fisher-Shannon plane ( $I-J$ ) in position space (red triangles) for the IRC of the abstraction reaction on a double logarithmic scale. The lower bound is depicted by the blue line (Equation (11) with $n=3$ ).


Figure 9. Fisher-Shannon plane ( $I-J$ ) in momentum space (red triangles) for the IRC of the abstraction reaction on a double logarithmic scale. The lower bound is depicted by the blue line (Equation (11) with $n=3$ ).

The complementarity of the conjugated spaces, $r$ and $p$, in the course of the chemical reaction can be analysed from Figures 6 and 7. The general observation from these figures is that they reflect opposite behavior, i.e. for each of the most representative regions in the reaction ( $\mathrm{R} / \mathrm{P}, \mathrm{BCER}$, etc.) we observe states of maximum delocalizability in position space (Figure 6) corresponding to minimum uniformity in momentum space (Figure 7).

Figures 8 and 9 show plots (on a doublelogarithmic scale) of the $I_{r}-J_{r}$ and $I_{p}-J_{p}$ planes for the chemical reaction. At this point it is worth mentioning that there is a rigorous lower bound to the associated $C(F S)$ complexity, given in Equation (6), which is $C(F S)=I \cdot J \geq 3$ for both spaces.

In position space, Figure 8 indicates a division of the $I_{r}-J_{r}$ plane into two regions where the straight line $I J=3$ (drawn in the plane on a logarithmic scale) divides it into an 'allowed' (upper) and a 'forbidden' (lower) part. In position space (Figure 8) the R/P are characterized by maximum values for order and uncertainty. As the reaction proceeds, disorder increases at the expense of decreasing uncertainty from the $\mathrm{R} / \mathrm{P}$ to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions. The latter has highly localized densities, and disorder maintains its decreasing path at the expense of augmenting its uncertainty from the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ to the TS. It is worth mentioning that the BCER are not present in this plane. Chemically, along the IRC we observe monotonically decreasing behavior for Fisher information, although the structural changes for the position space densities diminish more significantly from the $\mathrm{R} / \mathrm{P}$ to the $B-B / F$ regions than the changes from $B-B / F$ to the TS. These changes are accompanied by an increase in structural localizability from the $\mathrm{R} / \mathrm{P}$ to the $B-B / F$ regions and, from this point to the TS, the opposite is observed, i.e. the TS has a more delocalized position space density.

Figure 9 depicts the corresponding measures for the $I_{p}-J_{p}$ plane and some general aspects can be noted: (i) remarkable isocomplex behavior (linearity), which is divided into two regions, from $\mathrm{R} / \mathrm{P}$ to BCER, and from BCER to the TS, and (ii) the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region is missing from this plane. In addition, we note from Figure 9 that the $\mathrm{R} / \mathrm{P}$ are characterized by maximum structural order in momentum space that diminishes up to the BCER, which has maximum disorder and maximum uncertainty. This behavior then reverts and the uncertainty as well as the disorder of the process from BCER up to the TS decrease. This last point possesses the minimum global uncertainty. From a chemical point of view, the process behaves in such a way that, from R/P to BCER, energy is accumulated by gradually delocalizing the momentum space densities. The opposite is observed from BCER to the TS regions, where energy is released in order to achieve the bond cleavage and the spin coupling processes as discussed above.

For the $I-J$ plane the behavior of the conjugated spaces in the course of the reaction appears to be more complex, in that not all the regions are fully caracterized in both spaces as commented on above. In position space (Figure 8), localizability increases from the $\mathrm{R} / \mathrm{P}$ to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions, whereas in momentum space (Figure 9) it diminishes from R/P to BCER. The opposite is observed from the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ to the TS in that the localizability diminishes in position space, whereas in momentum space it increases from the BCER to the TS.

### 3.3. Complexities

In the search for joint patterns of uniformitylocalizability through $C(L M C)$ and disorder- localizability through $C(F S)$ we show in Figures 10 and 11 the values for these complexity measures in position and momentum space, respectively. The general observation from Figure 10 is that both complexity measures behave similarly in position space. It can be observed that the $\mathrm{R} / \mathrm{P}$ regions have maximum complexity and, as the reaction evolves, both complexities diminish up to the TS, which has the minimum complexity value at the IRC. It is worth noting that both measures fail to detect the BCER and $B-B / F$ regions.

In momentum space the $C(L M C)$ measure looks very much like that in position space, i.e. maximum values for the $\mathrm{R} / \mathrm{P}$ regions and minimum for the TS , and failing to detect the BCER and $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions. The situation for the Fisher-Shannon measure is different. As the reaction evolves, complexity decreases from maximum values at the $\mathrm{R} / \mathrm{P}$ regions up to the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions, which have the minimum values, and the complexity increases so as to reach the TS. It is interesting to note that whereas the joint measure for $I_{p}$ and $J_{p}$ is able to detect the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ regions, neither of these two separate measures reveal it, as observed from Figures 1 and 3.

At this stage of the analysis we have managed with valuable data to establish a relationship between the information-theoretical features of the reaction studied in the present work and that of the total energy at its IRC profile. Thus, Figures $12-14$ depict the $C(L M C)$ and $C(F S)$ complexity values as a function of energy for the conjugated spaces and for the product space, respectively.

Figure 12 shows monotonic decreasing behavior for both complexity measures (in position space) versus energy, i.e. as the reaction evolves the complexity diminishes at the expense of augmenting the energy up to the TS, which is shown at the bottom right corner of the figure. Further, our results show fairly linear behavior along the IRC except for the region in the vicinity of the R/P. Hence, we note that, in position space, the energy profile of the abstraction reaction bears the joint features of $C(L M C)$ (uniformitylocalizability) and $C(F S)$ (disorder-localizability).

Also, the information-theoretic features of uniformity-localizability in momentum space shown by $C(L M C)$ in Figure 13 appear to describe the behavior of the total energy to a very good extent, except for a small region close to the $\mathrm{R} / \mathrm{P}$. In contrast, the $C(F S)$ measure behaves in a more complex manner with respect to the energy, and does not reflect the


Figure 10. $C(L M C)$ (red triangles) and $C(F S)$ (blue circles) in position space for the IRC of the abstraction reaction.


Figure 11. $C(L M C)$ (red triangles) and $C(F S)$ (blue circles) in momentum space for the IRC of the abstraction reaction.


Figure 12. $C(L M C)$ (red triangles) and $C(F S)$ (blue circles) in position space as a function of the total energy of the abstraction reaction.


Figure 13. $C(L M C)$ (red triangles) and $C(F S)$ (blue circles) in momentum space as a function of the total energy of the abstraction reaction.


Figure 14. $C(L M C)$ (red triangles) and $C(F S)$ (blue circles) in the product space $r-p$ as a function of the total energy of the abstraction reaction.
$B-B / F$ aspects of the reaction, which are indeed revealed by $C(F S)$.

It is interesting to collect together all the information-theoretical features analysed in this work through the complexities $C(L M C)$ and $C(F S)$ in product space to obtain a description of the energy profile for the abstraction reaction. As can be seen from Figure 14, the energy profile exhibits features of uniformity-localizability and disorderlocalizability in the joint space. Therefore, it seems feasible that most of the chemical features of interest, namely the BCER and the $B-B / F$, are lost in the energy profile. As observed above, these features are revealed only when the particular informationtheoretical aspects of uniformity, localizability and disorder are considered.

## 4. Conclusions

In this work, we have investigated the complexity of the hydrogenic abstraction reaction by means of information functionals $D, L, I$ and $J$ and joint information-theoretic measures, i.e. the $I-D, D-L$ and $I-J$ planes and the Fisher-Shannon and LMC shape complexities.

The analysis of the information-theoretical functionals of the one-particle density was performed in position $(r)$ and momentum ( $p$ ) space. These measures were found to reveal all the chemically significant aspects of the course of the reaction, i.e. the reactant/ product region, the bond-cleavage energy region, the bond-breaking/forming region and the transition state. In addition, the information-theoretical concepts of uniformity, disorder and localizability were used to reveal the chemical phenomena of energy accumulation/release and to identify the mechanisms for bond forming and spin coupling.

In addition, the Fisher-disequilibrium ( $I-D$ ), the $L M C(D-L)$ and the Fisher-Shannon $(I-J)$ planes were studied to identify informational aspects of disorderuniformity, uniformity-localizability and disorderlocalizability. The analysis revealed that all of the chemically significant regions can be identified from most information-theoretical planes as well as the energetic course of the reaction. It was found that although some planes are linear to the bounds of isocomplexity, the behavior of the reaction presents more complex patterns for some regions.

The complementarity of the conjugated spaces is manifest in a complex manner, i.e. for the planes we observe opposite behavior for some regions located between the chemically significant zones of the reaction, i.e. the $\mathrm{R} / \mathrm{P}, \mathrm{BCER}, \mathrm{B}-\mathrm{B} / \mathrm{F}$ and the TS. However, these features are not completely characterized in the $I-J$ planes (nor in the complexities) so as to clearly observe opposite behavior of the conjugated spaces. On the other hand, it may be noted that position space information measures and complexities are closely related to the structural changes of the reactive complex along the IRC (e.g., B-B/F), whereas those in momentum space are associated with kinetic energetic changes (e.g., BCER).

It is important to mention that the local behavior of the Fisher information measure, as calculated by $C(F S)$ in momentum space, is the only complexity measure that can describe the $\mathrm{B}-\mathrm{B} / \mathrm{F}$ region, as compared with the other global measures which only reveal the R/P and TS.

According to the analysis of the complexities we note that, in position as well as in the joint space, the energy profile of the abstraction reaction bears the
joint information-theoretical features of uniformitylocalizability through $C(L M C)$ and disorder-localizability through $C(F S)$. Finally, it is feasible that most of the chemical features of interest, namely the BCER and the $\mathrm{B}-\mathrm{B} / \mathrm{F}$, are lost in the energy profile; these features are revealed solely when the particular information-theoretical aspects of $D, L$ or $J$, and $I$ are considered.

The results of this study indicate that further investigations are necessary in order to improve our understanding of the complexity of chemical reactions along the lines of analysing different reaction mechanisms, other information functionals, more intrincate aspects of the energy profile, etc. We believe that the strategy followed in this study might be useful in more complex cases by describing the phenomenological behavior of the chemical probe concerning local and global features by use of three key information measures, i.e. Shannon, Fisher and disequilibrium. The information planes and complexity measures of the process might then be resolved in a feasible manner.

## Acknowledgements

We wish to thank José María Pérez-Jordá and Miroslav Kohout for kindly providing their numerical codes. R.O.E. wishes to thank Juan Carlos Angulo and Jesús SánchezDehesa for their kind hospitality during his sabbatical stay at the Departamento de Física Atómica, Molecular y Nuclear and the Instituto Carlos I de Física Teórica y Computacional at the Universidad de Granada, Spain. We acknowledge financial support via Mexican grants from CONACyT, PIFI and PROMEP-SEP and Spanish grants MICINN projects FIS-2008-02380, FQM-4643 and P06-FQM-2445 of Junta de Andalucía. J.A., J.C.A., J.S.D., and R.O.E. belong to the Andalusian research group FQM-0207. R.O.E. wishes to acknowledge financial support from the Ministerio de Educación of Spain through grant SAB2009-0120. Allocation of supercomputing time from Laboratorio de Supercómputo y Visualización at UAM, Sección de Supercomputacion at CSIRC Universidad de Granada, and Departamento de Supercómputo at DGSCA-UNAM is gratefully acknowledged.

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