

Existence and stability of periodic solutions of the relativistic oscillator

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Introduction

Classical pendulum equation:

$$x'' + kx' + a \sin x = p(t) \equiv \bar{p} + \tilde{p}(t) \quad (1)$$

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Questions:

- ▶ **Degeneracy problem:** is $I_{\tilde{p}}$ nondegenerate?
- ▶ If $k = 0$, it is known that $0 \in I_{\tilde{p}}$ (Hamel, 1922). Is this true for $k > 0$?

Introduction

Theorem (Ortega's counterexample)

Given positive constants a, k and T , there exists $p \in \tilde{\mathcal{C}}_T$ such that the equation (1) with $\bar{p} = 0$ has no T -periodic solutions.

Introduction

Relativistic pendulum equation:

$$\phi(x')' + kx' + a \sin x = p(t) \equiv \bar{p} + \tilde{p}(t) \quad (2)$$

where $\phi :] - c, c[\rightarrow \mathbb{R}$ is given by

$$\phi(u) = \frac{u}{\sqrt{1 - \frac{u^2}{c^2}}}.$$

$c \equiv$ Speed of light in the vacuum

Main result

$$(\phi(x'))' + h(x)x' + g(x) = p(t) \quad (3)$$

where h, g are continuous functions, p is continuous and T -periodic and $\phi :] - c, c[\rightarrow \mathbb{R}$ is an increasing homeomorphism with $\phi(0) = 0$.

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Theorem (Bereanu-Mawhin, JDE, 2007)

If $g : \mathbb{R} \rightarrow \mathbb{R}$ is strictly monotone then $I_{\tilde{p}} = \text{Range } g$ for any $\tilde{p} \in \tilde{C}_T$.

Main result

Theorem

Let us assume that there exist real numbers $R_1 < R_2$ such that

(H1) $g \in C^1([R_1, R_2])$ and $g'(x) > 0$ for all $x \in]R_1, R_2[$.

(H2) $\frac{cT}{\sqrt{3}} \max_{x \in [R_1, R_2]} |g'(x)| \leq |g(R_2) - g(R_1)|$.

Then, for every $\tilde{p} \in \tilde{\mathcal{C}}_T$, the set $I_{\tilde{p}}$ of eq. (3) contains the closed interval

$$\left[g(R_1) + \frac{cT}{2\sqrt{3}} \max_{x \in [R_1, R_2]} |g'(x)|, g(R_2) - \frac{cT}{2\sqrt{3}} \max_{x \in [R_1, R_2]} |g'(x)| \right]$$

Moreover, if \bar{p} belongs to the previous interval, then the corresponding T -periodic solution belongs to $[R_1, R_2]$.

Main result

Theorem

Let us assume that there exist real numbers $R_1 < R_2$ such that

(H1) $g \in C^1([R_1, R_2])$ and $g'(x) < 0$ for all $x \in]R_1, R_2[$.

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Then, for every $\tilde{p} \in \tilde{\mathcal{C}}_T$, the set $I_{\tilde{p}}$ of eq. (3) contains the closed interval

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Moreover, if \bar{p} belongs to the previous interval, then the corresponding T -periodic solution belongs to $[R_1, R_2]$.

Corollary for the relativistic pendulum

Corollary

Let us assume that $cT \leq 2\sqrt{3}$. Then, for every

$$|\bar{p}| \leq a \left(1 - \frac{cT}{2\sqrt{3}} \right),$$

eq. (2) possesses two different solutions x_1, x_2 which verify $-\frac{\pi}{2} \leq x_1 \leq \frac{\pi}{2} \leq x_2 \leq \frac{3\pi}{2}$.

Proof.

Step 1: change of variables $y = g(x) - \bar{p}$

$$\phi \left(\frac{y'}{g'(g^{-1}(y + \bar{p}))} \right)' + h(g^{-1}(y + \bar{p})) \frac{y'}{g'(g^{-1}(y + \bar{p}))} + y(t) = \tilde{p}(t). \quad (4)$$

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Step 2: Formulation of the fixed point problem.

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Step 2: Formulation of the fixed point problem.

Integrating,

$$\phi \left(\frac{y'}{g'(g^{-1}(y + \bar{p}))} \right) = \int_0^t (p(s) - y(s)) ds - H(g^{-1}(y + \bar{p})) + C,$$

where H is a primitive of the function h and C is a constant to be fixed later.

If

$$F[y](t) = \int_0^t (p(s) - y(s)) ds - H(g^{-1}(y + \bar{p})).$$

then

$$y(t) = \int_0^t g'(g^{-1}(y + \bar{p})) \phi^{-1} (F[y](t) + C) ds + D.$$

Proof.

Step 2: Formulation of the fixed point problem.

Lemma

For any $y \in \tilde{\mathcal{C}}_T$, there exists a unique choice of C_y, D_y such that

$$\mathcal{T}[y](t) \equiv \int_0^t g'(g^{-1}(y + \bar{p}))\phi^{-1} (F[y](t) + C_y) ds + D_y \in \tilde{\mathcal{C}}_T. \quad (5)$$

The functional $\mathcal{T} : \tilde{\mathcal{C}}_T \rightarrow \tilde{\mathcal{C}}_T$ is continuous and compact.

Proof.

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$$K = \{y \in \tilde{C}_T : y(t) \in [g(R_1) - \bar{p}, g(R_2) - \bar{p}]\}.$$

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Step 3: Application of the Schauder fixed point problem.

$$K = \{y \in \tilde{C}_T : y(t) \in [g(R_1) - \bar{p}, g(R_2) - \bar{p}]\}.$$

$$\|\mathcal{T}[y]\|_\infty \leq \frac{T}{2\sqrt{3}} \|\mathcal{T}[y]'\|_\infty < \frac{cT}{2\sqrt{3}} \max_{x \in [R_1, R_2]} |g'(x)|.$$

By using condition (H2), $\mathcal{T}[y] \in K$.

Open problems

Open problems

- ▶ Prove or disprove existence for $T \gg \cdot$.
- ▶ Stability

Stability

The relativistic (an)harmonic oscillator

$$\left(\frac{mx'}{\sqrt{1 - \frac{x'^2}{c^2}}} \right)' + kx = -F_0 \cos \omega t \quad (6)$$

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Chaotic behavior shown numerically in

- ▶ J.H. Kim, H.W. Lee, Relativistic chaos in the driven harmonic oscillator, Physical Rev. E, **51** n.2 (1995), pp. 1579-1581.

Stability: Main result

Theorem

Assume that the parameters m , k , ω and F_0 satisfy the following conditions

$$(H1) \quad \frac{k}{m} < \frac{\omega^2}{16},$$

$$(H2) \quad F_0 < \frac{1}{4} mc\omega$$

$$(H3) \quad \frac{(mc\omega)^{19}[(mc\omega)^2 - 16F_0^2]}{120\pi F_0^2 (c^2 m^2 \omega^2 + 4F_0^2)^{19/2}} \sin\left(\frac{6\pi\omega^{1/2} mc^{3/2} k^{1/2}}{(c^2 m^2 \omega^2 + 4F_0^2)^{3/4}}\right) > 1,$$

then the driven relativistic harmonic oscillator (6) has a unique stable $2\pi/\omega$ -periodic solution with a twist dynamics around it (\implies Lyapunov stability + generically KAM dynamics).

Stability: Corollaries

Corollary

Fixed ω, k, F_0, c in (6), there exists $M_0 \equiv M_0(\omega, k, F_0, c) > 0$ such that if $m > M_0$ then the conclusion of Theorem 7 holds.

Corollary

Fixed m, k, F_0, c in (6), there exists a critical frequency $\omega_0 \equiv \omega_0(m, k, F_0, c)$ such that if $\omega > \omega_0$ then the conclusion of Theorem 7 holds.

Existence: an equivalent Newtonian oscillator

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(6) can be written as

$$x' = \frac{c^2 y}{\sqrt{y^2 c^2 + m^2 c^4}}, \quad y' = -kx - F_0 \cos \omega t. \quad (7)$$

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By deriving in the second equation, this system is equivalent to the second order equation

$$y'' + f(y) = \omega F_0 \sin \omega t, \quad (8)$$

with

$$f(y) = \frac{kc^2 y}{\sqrt{y^2 c^2 + m^2 c^4}}. \quad (9)$$

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Proposition

Let us assume that $4k < m\omega^2$. Then, the equation (8) has a T -periodic solution φ such that

$$-\frac{F_0}{\omega}(1 + \sin \omega t) < \varphi(t) < \frac{F_0}{\omega}(1 - \sin \omega t) \quad (10)$$

for all t .

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for all t .

Proof.

- ▶ $\alpha(t) = \frac{F_0}{\omega}(1 - \sin \omega t)$ lower solution
- ▶ $\beta(t) = -\frac{F_0}{\omega}(1 + \sin \omega t)$ upper solution
- ▶ $\beta(t) < \alpha(t)$
- ▶ $|f'(s)| \leq \frac{k}{m} < \left(\frac{\pi}{T}\right)^2$

Uniqueness and linear stability.

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Proposition

Assume that

$$4k < m\omega^2 \tag{11}$$

Then, (8) has a unique a T -periodic solution φ which is elliptic.

Lyapunov stability.

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The periodic solution $\varphi(t)$ of a general equation $y'' + g(t, y) = 0$ is translated to the origin by the canonical change $x = y - \varphi(t)$. For the equivalent equation

$$x'' + g(t, x + \varphi(t)) - g(t, \varphi(t)) = 0, \quad (12)$$

the equilibrium $x \equiv 0$ is a solution. By developing the nonlinearity up to the third order,

$$x'' + a(t)x + b(t)x^2 + c(t)x^3 + R(t, x) = 0. \quad (13)$$

where

$$a(t) = g_x(t, \varphi(t)), b(t) = g_{xx}(t, \varphi(t))/2, c(t) = g_{xxx}(t, \varphi(t))/6.$$

Definition

We say that the equilibrium $x \equiv 0$ of (13) is of twist type if it is elliptic, not strongly resonant and the associated first twist coefficient $\beta \neq 0$.

Lyapunov stability.

Theorem

Assume that for the equation (13) the following conditions holds:

- i) $-c(t) > c_* > 0, \forall t \in \mathbb{R}$
- ii) $|b(t)| \leq b^*, \forall t \in \mathbb{R}$
- iii) $0 < \sigma_1^2 < a(t) < \sigma_2^2 \leq (\frac{\pi}{2T})^2, \forall t \in \mathbb{R},$
- iv) $c_* > \frac{10T\sigma_2^7}{3 \sin(\frac{3T}{2}\sigma_1)\sigma_1^8} b^{*2}$ (twist condition)

Then the equilibrium $x \equiv 0$ of (13) is of twist type.