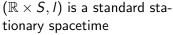
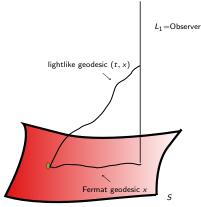
On the interplay between Lorentzian Causality and Finsler metrics of Randers type

Erasmo Caponio, Miguel Angel Javaloyes and Miguel Sánchez

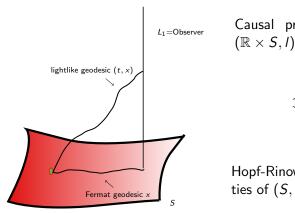
Universidad de Granada

International congress in Lorentzian geometry Martina Franca, July 8-11 (2009)





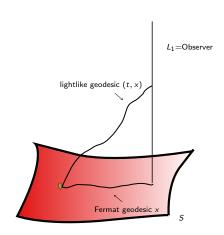
S is naturally endowed with a Randers metric F called the Fermat metric



Causal properties of $(\mathbb{R} \times S, I)$

 \updownarrow

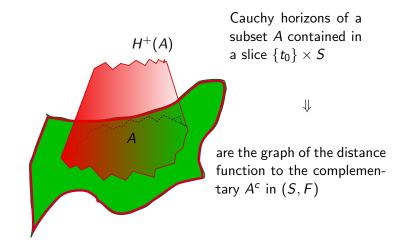
Hopf-Rinow properties of (S, F)



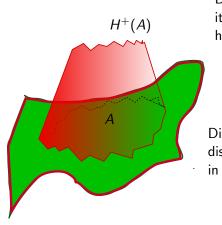
Global hyperbolicity of $(\mathbb{R} \times S, I)$

1

 $\bar{B}^+(p,r)\cap \bar{B}^-(p,r)$ compact $\forall p \in S$ and $\forall r > 0$ in (S,F)



◆□ > ◆□ > ◆□ > ◆□ > □ □



Differential properities of the Cauchy horizons in $(\mathbb{R} \times S, I)$



Differential properties of the distance function to a subset in (S, F)

• Preliminaries:

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- Second application: equivalence of differentiability of Cauchy horizons and the distance function to a subset.

Causal properties classify spacetimes depending on the behaviour of causal cones. A spacetime is:



Causally simple



Causally continuous



Stably causal



Strongly causal



Distinguishing



Causal



Chronological



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Globally hyperbolic



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Non-totally vicious

E. Caponio,

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- Globally hyperbolic if it admits a Cauchy hypersurface (a subset S that meets exactly once every inextendible timelike curve)



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M. A. J. AND M. SÁNCHEZ, A note on the existence of standard splittings for conformally stationary spacetimes,

Classical Quantum Gravity, 25 (2008), pp. 168001, 7.



Theorem (M. A. J.- M. Sánchez)

If a stationary spacetime L is distinguishing and the timelike Killing field is complete, then it is causally continuous and standard



Causally simple



Causally continuous

Stably causal

Strongly causal



Distinguishing



Causal



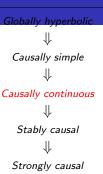
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Sketch of the proof:



Distinguishing

Causal

Chronological

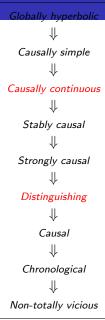
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6 / 26

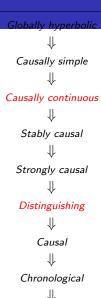
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Causally simple Causally continuous Stably causal Strongly causal Distinguishing Causal

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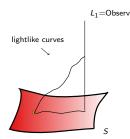
Chronological



Fermat principle in standard stationary spacetimes

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 Relativistic Fermat Principle: lightlike pregeodesics are critical points of the arrival time function corresponding to an observer in a suitable class of lightlike curves



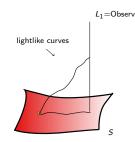


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$$t(b) \! = \! t(a) + \! \int_a^b \! \left(\tfrac{1}{\beta} g_0(\dot{x},\!\delta) \! + \! \sqrt{\tfrac{1}{\beta} g_0(\dot{x},\!\dot{x}) \! + \! \tfrac{1}{\beta^2} g_0(\dot{x},\!\delta)^2} \right) \! \mathrm{d}s.$$





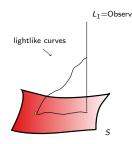
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PIERRE DE FERMAT (1601-1665)

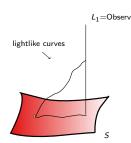
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- ullet Let us define the Fermat (Finslerian) metric in S as

$$F(x,v) = \frac{1}{\beta} g_0(v,\delta) + \sqrt{\frac{1}{\beta} g_0(v,v) + \frac{1}{\beta^2} g_0(v,\delta)^2},$$





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A curve $s \to \gamma(s) = (s, x(s))$ is a lightlike pregeodesic of $(\mathbb{R} \times S, g)$ iff $s \to x(s)$ is a Fermat geodesic with unit speed.

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Einstein Cross



Gravitational lensing

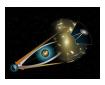
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 - Existence of t-periodic lightlike geodesics is equivalent to existence of Fermat closed geodesics (Biliotti-M.A.J. to appear in Houston J. Math.)



Einstein Cross



Gravitational lensing

Randers metrics

Randers metrics

• Randers metrics in a manifold M is a function $R: TM \to \mathbb{R}$ defined as:

$$R(x, v) = \sqrt{h(v, v)} + \omega_x[v]$$

where h is Riemannian and ω a 1-form with $\|\omega_x\|_h < 1 \ \forall x \in M$, are basic examples of non-reversible Finsler metrics: $R(x, -v) \neq R(x, v)$.

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 Named after the norwegian physicist Gunnar Randers (1914-1992):

Randers, G.: On an asymmetrical metric in GINNAR RANDERS WITH ALBERT EINSTEIN the fourspace of General Relativity. Phys. Rev. (2) **59**, 195–199 (1941)





Main reference:



Bao, D., Chern, S.S., Shen, Z.: An Introduction to Riemann-Finsler geometry.

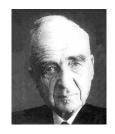
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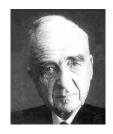
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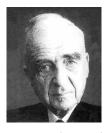
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- **3** Fiberwise strictly convex square: $g_{ij}(x,y) = \left[\frac{1}{2} \frac{\partial^2 (F^2)}{\partial y^i \partial y^j}(x,y)\right]$ is positively defined.



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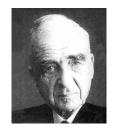
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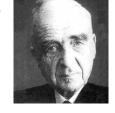
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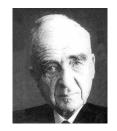
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- Define the symmetrized distance

$$\mathrm{d}_{s}(p,q) = \frac{1}{2}(\mathrm{d}(p,q) + \mathrm{d}(q,p))$$

and
$$B_s(x,r) = \{ p \in S : d_s(x,p) < r \}$$

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- Define the symmetrized distance

$$\mathrm{d}_s(p,q) = \frac{1}{2}(\mathrm{d}(p,q) + \mathrm{d}(q,p))$$

and
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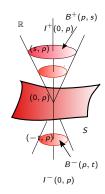
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Globally hyperbolic



Causally simple



Causally continuous



Stably causal



Strongly causal



Distinguishing



Causal



Chronological



Non-totally vicious

Theorem

E. Caponio,

Theorem

Let $(\mathbb{R} \times S, g)$ be a standard stationary spacetime. Then $(\mathbb{R} \times S, g)$ is causally continuous and



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Causality through the Fermat metric

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- (c) a slice $\{t_0\} \times S$, $t_0 \in \mathbb{R}$, is a Cauchy hypersurface if and only if the Fermat metric F on S is forward and backward complete.

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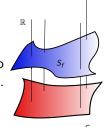


Non-totally vicious

Randers metrics with the same geodesics

Randers metrics with the same geodesics

• Let R and R' be Randers metrics. They are associated to R' the same stationary spacetime if and only if R' = R + df.



$$S_f = \{ (f(x), x) : x \in S \}$$
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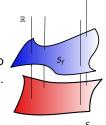
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Randers metrics with the same geodesics

- Let R and R' be Randers metrics. They are associated to the same stationary spacetime if and only if R' = R + df.
- Moreover, if $\mathbb{R} \times S$ is the splitting associated to R, the splitting associated to R' is $\mathbb{R} \times S_f$, where

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Theorem (Accurate Hopf-Rinow for Randers metrics)



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(A) the intersection $\bar{B}^+(x,r) \cap \bar{B}^-(x,r)$ of (S,R) is compact for every r > 0 and $x \in S$



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In such a case, (S, R) is convex.



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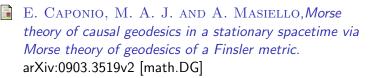


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 - E. CAPONIO, M. A. J. AND A. MASIELLO, Morse theory of causal geodesics in a stationary spacetime via Morse theory of geodesics of a Finsler metric. arXiv:0903.3519v2 [math.DG]
- As an application we obtain Morse theory for lightlike geodesics and timelike geodesics with fixed proper time from a point to a vertical line.

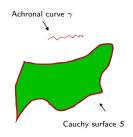
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D^{\pm}(A) = \{ p \in M : \text{ every past (resp. future)}  inextendible causal curve through p meets A \}
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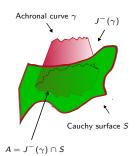
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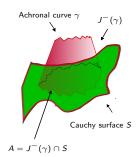
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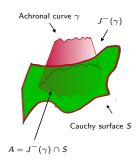
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• the future (resp. past) Cauchy horizon is

$$H^{\pm}(A) = \{ p \in D^{\pm}(A) : I^{\pm}(p) \text{ does not meet } D^{\pm}(A) \}$$

Theorem

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$$\inf_{x \notin A} d(y,x) = d(y,A^c)$$

Theorem

Let $(\mathbb{R} \times S, g)$ be a standard stationary spacetime such that $\{t_0\} \times S$ is Cauchy, and $A_{t_0} = \{t_0\} \times A$. Then

$$D^{+}(A_{t_0}) = \{(t, y) : d(x, y) > t - t_0 \ \forall x \notin A \text{ and } t \geq t_0\},$$

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Cauchy horizons can be seen as the graph of the distance function to a subset!!!!

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Li-Nirenberg theorem

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Yan'yan Li and Louis Nirenber

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Theorem (Li-Nirenberg)

The function $\partial\Omega\ni y\to\min(N,\ell(y))\in\mathbb{R}^+$ is Lipschitz-continuous on any compact subset. As a consequence $\mathfrak{h}^{n-1}(\Sigma\cap B)<+\infty$, being B bounded.



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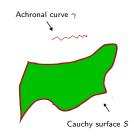
Y. LI AND L. NIRENBERG, The distance function to the boundary, Finsler geometry, and the singular set of viscosity solutions of some Hamilton-Jacobi equations, Comm. Pure Appl. Math., (2005).



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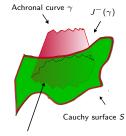


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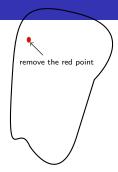




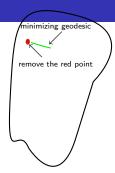
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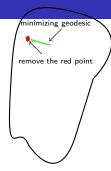
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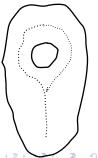
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- Y. LI AND L. NIRENBERG, The distance function to the boundary, Finsler geometry, and the singular set of viscosity solutions of some Hamilton-Jacobi equations, Comm. Pure Appl. Math., (2005).





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- There are several results for the differentiability of future horizons:
- J. K. Beem and A. Królak, Cauchy horizon end points and differentiability,
 - J. Math. Phys., 39 (1998), pp. 6001–6010.
- P. T. CHRUŚCIEL, J. H. G. Fu, G. J. GALLOWAY, AND R. HOWARD, On fine differentiability properties of horizons and applications to Riemannian geometry,
 - J. Geom. Phys., 41 (2002), pp. 1–12.

Cut loci of Randers metrics via Cauchy horizons

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Corollary

The n-dimensional Haussdorf measure of Cut_C is zero.

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 - G. W. GIBBONS, C. A. R. HERDEIRO, C. M. WARNICK, M. C. WERNER, Stationary Metrics and Optical Zermelo-Randers-Finsler Geometry.

Phys.Rev.D79: 044022,2009

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the authors show that Fermat metrics with constant flag curvature correspond with locally conformally flat stationary spacetimes, but the converse is not true.

(3) Which is the condition in the Fermat metric that characterizes conformally flatness for the stationary spacetime?

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