



Constructing matrix-valued coordinate transformations from scratch

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In collaboration with: J. Adam, W. Troost and W. Van Herck (K.U.Leuven)

References: [arXiv:0712.0918](https://arxiv.org/abs/0712.0918) [hep-th].

Motivation

A lot has been learned about the dynamics of **multiple D-branes** in the last past years:

- $U(1)^N \rightarrow U(N)$ symmetry enhancement
- effective actions describing the non-Abelian dynamics
- many applications of non-Abelian effects in modern string theory

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Question: Is the action **invariant under background gauge transformations** ?

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Question: Is the action **invariant under background general coord transf** ?

- **Worldvolume coordinates** : Abelian theory \longrightarrow trivial
- **Target space coordinates** : manifest as **adjoint scalars** X^μ in worldvolume
How do background diffeomorphisms affect matrix coordinates X^μ ?

Outlook

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- Physics of multiple branes
- Problems with diffeomorphisms

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2. Matrix-valued differential geometry from scratch

- Matrix scalars
- Matrix contravariant vectors
- Matrix covariant vectors
- Matrix tensors and form fields
- Scalar product & problems

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3. Conclusions

1. Non-Abelian physics of multiple parallel branes

The physics of N **separated** parallel Dp -branes is very different from physics of N **coinciding** Dp -branes.

- **separated**: \rightarrow **Abelian** theory
- **coinciding**: \rightarrow **non-Abelian** theory

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This non-Abelian character has numerous manifestations in modern string theory:

- Dielectric effect [Myers]
- Gravity duals of (confining) gauge theories [Polchinsky, Strassler]
- Enhançons [Johnson]
- Matrix models in non-trivial backgrounds [Berenstein, Maldacena, Nastase]
- Microscopic description of giant gravitons [B.J., Lozano, Rodríguez]
- ...

Difference in degrees of freedom

- **separated:** N $U(1)$ vector fields V_a^I , $(9 - p)N$ scalars X^{iI}
→ **Abelian** worldvolume action
- **coinciding:** 1 $U(N)$ Yang-Mills vector V_a^I , N adjoint scalars X^{iI}
→ **non-Abelian** worldvolume action

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Extra degrees of freedom come from massless strings stretched between coinciding strings:

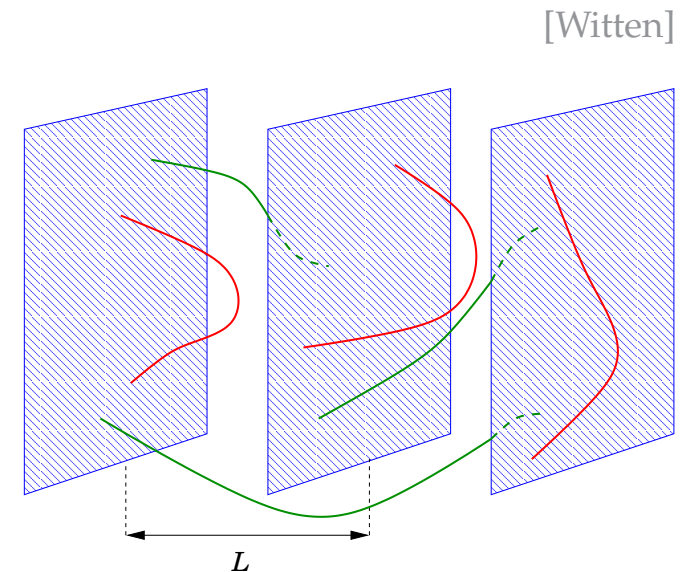
strings between same brane: $m \sim 0$

strings between different branes: $m \sim L$

As $L \rightarrow 0$:

⇒ $N + N(N - 1) = N^2$ degrees of freedom

⇒ $U(1)^N \rightarrow U(N)$ gauge enhancement



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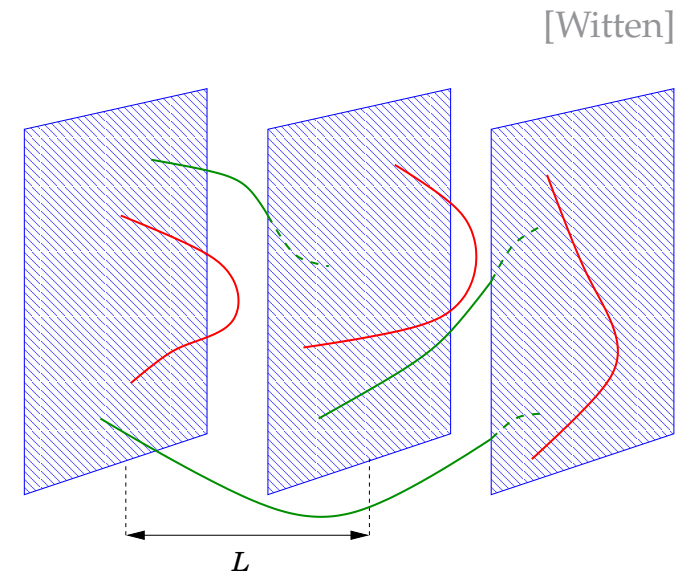
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Extra degrees of freedom and non-Abelian symmetries

⇒ non-Abelian worldvolume action



Role of the scalars

- Abelian: $X^{\mu I}(\sigma)$ indicate position of I th brane in direction x^μ
→ rearrange in diagonal matrix $X^\mu(\sigma) = \text{diag}\left(X^{\mu 1}(\sigma), \dots, X^{\mu N}(\sigma)\right)$

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- Background fields: $\Phi = \Phi(x) \longrightarrow \Phi = \Phi(X)$
→ $\Phi(X) = \sum_n \frac{1}{n!} \partial_{\mu_1} \dots \partial_{\mu_n} \Phi(x)|_{x^\lambda=0} X^{\mu_1} \dots X^{\mu_n}$ [Douglas][Garousi, Myers]
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- coordinate transf: $g'_{\alpha\beta} = \Lambda^\mu_\alpha \Lambda^\nu_\beta g_{\mu\nu} \longrightarrow X' = X'(X)$ such that

$$S[g, X] \equiv S[g', X']$$

Question: What is $X' = X'(X)$?

→ Problems when combining ordering and transitivity

[de Boer, Schalm]

Abelian composition:

$$y^\mu = x^\mu + a_{\nu\rho}^\mu x^\nu x^\rho + b_{\nu\rho\lambda}^\mu x^\nu x^\rho x^\lambda + \dots$$

$$z^\mu = y^\mu + \tilde{a}_{\nu\rho}^\mu y^\nu y^\rho + \tilde{b}_{\nu\rho\lambda}^\mu y^\nu y^\rho y^\lambda + \dots$$

$$= x^\mu + (a_{\nu\rho}^\mu + \tilde{a}_{\nu\rho}^\mu) x^\nu x^\rho + (b_{\nu\rho\lambda}^\mu + 2a_{\nu\sigma}^\mu \tilde{a}_{\rho\lambda}^\sigma + \tilde{b}_{\nu\rho\lambda}^\mu) x^\nu x^\rho x^\lambda + \dots$$

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Non-abelian composition:

$$Y^\mu = X^\mu + F_{\nu\rho}^\mu(a) X^\nu X^\rho + G_{\nu\rho\lambda}^\mu(a, b) X^\nu X^\rho X^\lambda + \dots$$

$$Z^\mu = Y^\mu + F_{\nu\rho}^\mu(\tilde{a}) Y^\nu Y^\rho + G_{\nu\rho\lambda}^\mu(\tilde{a}, \tilde{b}) Y^\nu Y^\rho Y^\lambda + \dots$$

$$= X^\mu + \left[F_{\nu\rho}^\mu(a) + F_{\nu\rho}^\mu(\tilde{a}) \right] X^\nu X^\rho$$

$$+ \left[G_{\nu\rho\lambda}^\mu(a, b) + F_{\nu\sigma}^\mu(\tilde{a}) F_{\rho\lambda}^\sigma(a) + F_{\sigma\lambda}^\mu(\tilde{a}) F_{\nu\rho}^\sigma(a) + \tilde{G}_{\nu\rho\lambda}^\mu(\tilde{a}, \tilde{b}) \right] X^\nu X^\rho X^\lambda + \dots$$

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→ $\nexists F_{\nu\rho}^\mu$ and $G_{\nu\rho\lambda}^\mu$ such that above is true

$$f(x) \longrightarrow F(X), \quad g(x) \longrightarrow G(X)$$

$$h(x) = (g \circ f)(x) \longrightarrow H(X) \neq (G \circ F)(X) \quad \text{for any ordering!}$$

**Abelian diffeomorphisms is not a subset
of Non-Abelian DIFFEOMORPHISMS**

[de Boer, Schalm]

No solution for $S[g, X] = S[g'(g), X'(X)] \dots$

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Two strategies:

- Look for **inclusion function** via Noether procedure
- Construct **matrix differential geometry** from scratch
→ **rest of the talk**

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Construct an algebra for matrix coord transf on its own,
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by analogy of Abelian counterparts

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Matrix function, defined by expansion series:

$$F(X) = \sum_{k=0}^{\infty} \frac{1}{k!} f_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_k}$$

where $f_{\lambda_1 \dots \lambda_k}$ complex coefficients without specific symmetries

→ sum, product and composition gives matrix function

2.1 Substitution operator and matrix scalar

Coordinate transformation:

$$\begin{array}{ll} \text{Abelian :} & x^\mu \rightarrow y^\mu = x^\mu - \xi^\mu(x) \quad \text{with} \quad \xi(x) = \left(\xi^0(x), \dots, \xi^{D-1}(x) \right) \\ \text{non - Abelian :} & X^\mu \rightarrow Y^\mu = X^\mu - \Xi^\mu(X) \quad \text{with} \quad \Xi(x) = \left(\Xi^0(x), \dots, \Xi^{D-1}(x) \right) \end{array}$$

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Abelian scalars:

$$f'(y) = f(x) \quad \iff \quad \delta_\xi f(x) = f'(x) - f(x)$$

To first order in ξ

$$f(x) = \sum_{k=0}^{\infty} \frac{1}{k!} \left(\partial_{\lambda_1} \dots \partial_{\lambda_k} f \right) (0) x^{\lambda_1} \dots x^{\lambda_k}$$

$$\begin{aligned} \delta_\xi f(x) &= \sum_{k=0}^{\infty} \frac{1}{k!} \left(\partial_{\lambda_1} \dots \partial_{\lambda_k} f \right) (0) \left(\xi^{\lambda_1}(x) x^{\lambda_2} \dots x^{\lambda_k} + \dots + x^{\lambda_1} \dots x^{\lambda_{k-1}} \xi^{\lambda_k}(x) \right) \\ &= \xi^\rho \partial_\rho f(x) \end{aligned}$$

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where we defined $\overline{\Xi^\rho \partial_\rho}$ as:

“take each X^μ in $F(X)$ in turn and substitute by $\Xi^\mu(X)$ and sum over everything”

Hence **matrix scalars** are defined by

$$\delta_\Xi F(X) = \overline{\Xi^\rho \partial_\rho} F(X)$$

$$[\delta_1, \delta_2]F(X) = \overline{\Xi^\lambda \partial_\lambda} \left(\overline{\Sigma^\rho \partial_\rho} F(X) \right) - \overline{\Sigma^\lambda \partial_\lambda} \left(\overline{\Xi^\rho \partial_\rho} F(X) \right)$$

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&\quad - \left(\overline{\Sigma^\lambda \partial_\lambda} \middle| \overline{\Xi^\rho \partial_\rho} \right) F(X) - \overline{\left(\overline{\Sigma^\lambda \partial_\lambda} \Xi^\rho \right)} \partial_\rho F(X)
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$$\Lambda^\rho = \overline{\Xi^\lambda \partial_\lambda} \Sigma^\rho - \overline{\Sigma^\lambda \partial_\lambda} \Xi^\rho$$

- $\delta_{\Xi} F(X) = \overline{\Xi^\rho \partial_\rho} F(X)$ forms an algebra
- Leibnitz rule: $\overline{\Xi^\rho \partial_\rho} (F \cdot G) = \overline{\Xi^\rho \partial_\rho} F \cdot G + F \cdot \overline{\Xi^\rho \partial_\rho} G$

2.2 Contravariant matrix vectors

Abelian:

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Commutator:

$$\begin{aligned} [\delta_1, \delta_2] A^\mu &= \overline{\Xi^\rho \partial_\rho} \left(\overline{\Sigma^\sigma \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \Sigma^\mu \right) - \overline{\Sigma^\rho \partial_\rho} \left(\overline{\Xi^\sigma \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \Xi^\mu \right) \\ &= \overline{\left(\overline{\Xi^\rho \partial_\rho} \Sigma^\sigma - \overline{\Sigma^\rho \partial_\rho} \Xi^\sigma \right) \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \left(\overline{\Xi^\sigma \partial_\sigma} \Sigma^\mu - \overline{\Sigma^\sigma \partial_\sigma} \Xi^\mu \right) \end{aligned}$$

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$$\delta_\Xi A^\mu = \overline{\Xi^\rho \partial_\rho} A^\mu - \overline{A^\rho \partial_\rho} \Xi^\mu$$

Commutator:

$$\begin{aligned} [\delta_1, \delta_2] A^\mu &= \overline{\Xi^\rho \partial_\rho} \left(\overline{\Sigma^\sigma \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \Sigma^\mu \right) - \overline{\Sigma^\rho \partial_\rho} \left(\overline{\Xi^\sigma \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \Xi^\mu \right) \\ &= \overline{\left(\overline{\Xi^\rho \partial_\rho} \Sigma^\sigma - \overline{\Sigma^\rho \partial_\rho} \Xi^\sigma \right) \partial_\sigma} A^\mu - \overline{A^\rho \partial_\rho} \left(\overline{\Xi^\sigma \partial_\sigma} \Sigma^\mu - \overline{\Sigma^\sigma \partial_\sigma} \Xi^\mu \right) \end{aligned}$$

Scalar multiplication:

$$\begin{aligned} \delta_\Xi (F \cdot A^\mu) &= \delta_\Xi F \cdot A^\mu + F \cdot \delta_\Xi A^\mu = \overline{\Xi^\rho \partial_\rho} (F \cdot A^\mu) - F \cdot \overline{A^\rho \partial_\rho} \Xi \\ &\neq \overline{\Xi^\rho \partial_\rho} (F \cdot A^\mu) - \overline{(F \cdot A^\rho) \partial_\rho} \Xi^\mu \end{aligned}$$

→ $A^\mu(X)$ do not form vector space!

→ Give up Leibnitz rule? Enhance/restrict definition of matrix vector?

2.3 Covariant matrix vectors

Abelian: $df = dx^\rho \partial_\rho f$

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Non-Abelian:

$$\begin{aligned} \underline{d}F &= \sum_{k=0}^{\infty} \frac{1}{k!} f_{\lambda_1 \dots \lambda_k} \left(\underline{d}X^{\lambda_1} X^{\lambda_2} \dots X^{\lambda_k} + \dots + X^{\lambda_1} \dots X^{\lambda_{k-1}} \underline{d}X^{\lambda_k} \right) \\ &\equiv \overline{\underline{d}X^\rho \partial_\rho} F \end{aligned}$$

where we defined $\overline{\underline{d}X^\rho \partial_\rho}$ as

“take each X^μ in $F(X)$ in turn and substitute by $\underline{d}X^\mu$ and sum over everything”

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Variation:

$$\delta_\Xi (\underline{d}F) = \underline{d}(\delta_\Xi F) = \underline{d}(\overline{\Xi^\mu \partial_\mu} F) = \overline{\underline{d}\Xi^\mu \partial_\mu} F + \overline{\Xi^\mu \partial_\mu} \underline{d}F$$

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→ possible to write in terms of $\underline{d}F$ only?
(necessary for defining one-forms)

$\overline{d}\Xi^\mu \partial_\mu F$ reads

“take each X^μ in $F(X)$ in turn and substitute by $\underline{d}\Xi^\mu$ and sum over everything”

$$\implies \overline{d}\Xi^\mu \partial_\mu F = \left(\overline{\underline{d}\Xi^\mu \underline{\partial}_\mu} \circ \underline{d} \right) F = \overline{\underline{d}\Xi^\mu \underline{\partial}_\mu} (\underline{d}F)$$

where we defined $\overline{\underline{d}\Xi^\mu \underline{\partial}_\mu}$ as

“take each $\underline{d}X^\mu$ in $\underline{d}F(X)$ in turn (leaving X^ρ 's untouched) and substitute by $\underline{d}\Xi^\mu$ and sum over everything”

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Hence the variation of $\underline{d}F$ is

$$\underline{d}F = \sum_{k=0}^{\infty} \frac{1}{k!} f_{\lambda_1 \dots \lambda_k} \left(\underline{d}X^{\lambda_1} X^{\lambda_2} \dots X^{\lambda_k} + \dots + X^{\lambda_1} \dots X^{\lambda_{k-1}} \underline{d}X^{\lambda_k} \right)$$

$$\delta_\Xi(\underline{d}F) = \overline{\underline{d}\Xi^\mu \underline{\partial}_\mu}(\underline{d}F) + \overline{\Xi^\mu \partial_\mu}(\underline{d}F)$$

→ Generalize to one-forms that are not dF

Matrix one-forms

$$B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} \underline{d}X^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k}$$

$$\delta_{\Xi}(B) = \overline{d\Xi^{\mu} \partial_{\mu}}(B) + \overline{\Xi^{\mu} \partial_{\mu}}(B)$$

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Comments

- $[\delta_1, \delta_2]B = \overline{d\Lambda^{\mu} \partial_{\mu}}(B) + \overline{\Lambda^{\mu} \partial_{\mu}}(B)$

$$\text{with } \Lambda^{\mu} = \overline{\Xi^{\rho} \partial_{\rho}} \Sigma^{\mu} - \overline{\Sigma^{\rho} \partial_{\rho}} \Xi^{\mu}$$

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with $\Lambda^{\mu} = \overline{\Xi^{\rho} \partial_{\rho}} \Sigma^{\mu} - \overline{\Sigma^{\rho} \partial_{\rho}} \Xi^{\mu}$
- Not decomposable in components $B \neq B_{\mu} \underline{d}X^{\mu}$
- B matrix one-form, F, F' matrix scalars:
 - $F \cdot B \cdot F'$ is **matrix one-form**
 - Matrix one-forms span **vector space**

2.4 Covariant matrix tensors

Abelian:
$$C = C_{\mu\nu} dx^\mu \otimes dx^\nu$$

→ Place of dx^μ and dx^ν is important

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$$C = \sum_{k=2}^{\infty} \sum_{\substack{i,j=1 \\ i \neq j}}^k c_{\mu_1 \dots \mu_k}^{(i,j)} X^{\mu_1} \dots \underline{d}_{(1)} X^{\mu_i} \dots \underline{d}_{(2)} X^{\mu_j} \dots X^{\mu_k}$$

where $\underline{d}_{(i)} X^\mu$ indicates i th factor of tensor space

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Transformation: $\delta_\Xi C = \overline{\Xi^\rho} \partial_\rho C + \underline{\overline{d}\Xi^\rho} \underline{\partial}_\rho C$

Cfr: $\delta C_{\mu\nu} = \xi^\rho \partial_\rho C_{\mu\nu} + C_{\rho\nu} \partial_\mu \xi^\rho + C_{\mu\rho} \partial_\nu \xi^\rho$

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Transformation: $\delta_\Xi C = \overline{\Xi^\rho \partial_\rho} C + \overline{d\Xi^\rho \underline{d}_\rho} C$

Cfr: $\delta C_{\mu\nu} = \xi^\rho \partial_\rho C_{\mu\nu} + C_{\rho\nu} \partial_\mu \xi^\rho + C_{\mu\rho} \partial_\nu \xi^\rho$

Anti-symmetry: $c_{\mu_1 \dots \mu_k}^{(i,j)} = -c_{\mu_1 \dots \mu_k}^{(j,i)}$

$$C = \sum_{k=2}^{\infty} \sum_{\substack{i,j=1 \\ i \neq j}}^k c_{\mu_1 \dots \mu_k}^{(i,j)} \left(X^{\mu_1} \dots \underline{d}_{(1)} X^{\mu_i} \dots \underline{d}_{(2)} X^{\mu_j} \dots X^{\mu_k} - X^{\mu_1} \dots \underline{d}_{(2)} X^{\mu_i} \dots \underline{d}_{(1)} X^{\mu_j} \dots X^{\mu_k} \right)$$

$$\begin{aligned}
\underline{d} B &\equiv \sum_{k=2}^{\infty} \left[b_{\mu_1 \dots \mu_k}^{(1)} \left(\underline{d}_{(1)} X^{\mu_1} \underline{d}_{(2)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_k} + \underline{d}_{(1)} X^{\mu_1} X^{\mu_2} \underline{d}_{(2)} X^{\mu_3} \dots X^{\mu_k} \right. \right. \\
&\quad \left. \left. + \dots + \underline{d}_{(1)} X^{\mu_1} X^{\mu_2} \dots X^{\mu_{k-1}} \underline{d}_{(2)} X^{\mu_k} \right) \right. \\
&\quad \left. - b_{\mu_1 \dots \mu_k}^{(1)} \left(\underline{d}_{(2)} X^{\mu_1} \underline{d}_{(1)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_k} + \underline{d}_{(2)} X^{\mu_1} X^{\mu_2} \underline{d}_{(1)} X^{\mu_3} \dots X^{\mu_k} \right. \right. \\
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&\quad \left. + \dots + \right. \\
&\quad \left. + b_{\mu_1 \dots \mu_k}^{(k)} \left(\underline{d}_{(2)} X^{\mu_1} X^{\mu_2} \dots X^{\mu_{k-1}} \underline{d}_{(1)} X^{\mu_k} + X^{\mu_1} \underline{d}_{(2)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_{k-1}} \underline{d}_{(1)} X^{\mu_k} \right. \right. \\
&\quad \left. \left. + \dots + X^{\mu_1} \dots X^{\mu_{k-2}} \underline{d}_{(2)} X^{\mu_{k-1}} \underline{d}_{(1)} X^{\mu_k} \right) \right. \\
&\quad \left. - b_{\mu_1 \dots \mu_k}^{(k)} \left(\underline{d}_{(1)} X^{\mu_1} X^{\mu_2} \dots X^{\mu_{k-1}} \underline{d}_{(2)} X^{\mu_k} + X^{\mu_1} \underline{d}_{(1)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_{k-1}} \underline{d}_{(2)} X^{\mu_k} \right. \right. \\
&\quad \left. \left. + \dots + X^{\mu_1} \dots X^{\mu_{k-2}} \underline{d}_{(1)} X^{\mu_{k-1}} \underline{d}_{(2)} X^{\mu_k} \right) \right]
\end{aligned}$$

$$\begin{aligned}
\underline{d}B &\equiv \sum_{k=2}^{\infty} \left[b_{\mu_1 \dots \mu_k}^{(1)} \left(\underline{d}_{(1)} X^{\mu_1} \underline{d}_{(2)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_k} + \underline{d}_{(1)} X^{\mu_1} X^{\mu_2} \underline{d}_{(2)} X^{\mu_3} \dots X^{\mu_k} \right. \right. \\
&\quad \left. \left. + \dots + \underline{d}_{(1)} X^{\mu_1} X^{\mu_2} \dots X^{\mu_{k-1}} \underline{d}_{(2)} X^{\mu_k} \right) \right. \\
&\quad \left. - b_{\mu_1 \dots \mu_k}^{(1)} \left(\underline{d}_{(2)} X^{\mu_1} \underline{d}_{(1)} X^{\mu_2} X^{\mu_3} \dots X^{\mu_k} + \underline{d}_{(2)} X^{\mu_1} X^{\mu_2} \underline{d}_{(1)} X^{\mu_3} \dots X^{\mu_k} \right. \right. \\
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\end{aligned}$$

NB: $B = \underline{d}F \iff \underline{d}B = 0$

$\longrightarrow \underline{d}(\underline{d}.) = 0$

2.5 Scalar product

$$F = \sum_{k=0}^{\infty} \frac{1}{k!} f_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_k} \quad (\text{scalar})$$

$$A^\mu = \sum_{k=0}^{\infty} \frac{1}{k!} a_{\lambda_1 \dots \lambda_k}^\mu X^{\lambda_1} \dots X^{\lambda_k} \quad (\text{contrav. vector})$$

$$B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} \underline{d} X^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k} \quad (\text{covar. vector})$$

→ can we define a scalar product between covar and contrav vectors?

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→ can we define a scalar product between covar and contrav vectors?

→ substitute $\underline{d}X^\lambda$ in B by A^λ :

$$A \cdot B = \overline{A^\mu \underline{\partial}_\mu} B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} A^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k} = F(X)$$

where $\overline{A^\mu \underline{\partial}_\mu}$ reads

“take each $\underline{d}X^\mu$ in B in turn and substitute by A^μ and sum over everything”

Transformation: use Leibnitz rule

$$\begin{aligned}\delta_{\Xi}(A \cdot B) &= \delta_{\Xi}A \cdot B + A \cdot \delta_{\Xi}B \\ &= \overline{(\Xi^{\rho} \partial_{\rho} A^{\mu} - A^{\rho} \partial_{\rho} \Xi^{\mu})} \underline{\partial}_{\mu} B + \overline{A^{\mu}} \underline{\partial}_{\mu} (\overline{\Xi^{\rho}} \partial_{\rho} B + \underline{d}\overline{\Xi^{\rho}} \underline{\partial}_{\rho} B)\end{aligned}$$

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$$\text{NB : } \overline{(A^{\rho} \partial_{\rho} \Xi^{\mu})} \underline{\partial}_{\mu} B = \overline{A^{\mu} \underline{\partial}_{\mu}} (\underline{d\Xi^{\rho} \underline{\partial}_{\rho}} B)$$

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$$= \overline{\Xi^{\rho} \partial_{\rho}} (A \cdot B) \quad \longrightarrow \quad A^{\mu} \text{ and } B \text{ are dual objects}$$

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Comments and problems

- Scalar transformation: $\overline{\Xi^{\rho} \partial_{\rho}} F = \overline{\Xi^{\rho}} \underline{\partial}_{\rho} \underline{d}F = \Xi \cdot \underline{d}F$

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Comments and problems

- Scalar transformation: $\overline{\Xi^{\rho} \partial_{\rho}} F = \overline{\Xi^{\rho}} \underline{\partial}_{\rho} \underline{d}F = \Xi \cdot \underline{d}F$

- Every form dual to $\overline{A^{\rho}} \underline{\partial}_{\rho}$ is of the form

$$B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} \underline{d}X^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k}$$

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 \delta_{\Xi}(A \cdot B) &= \delta_{\Xi}A \cdot B + A \cdot \delta_{\Xi}B \\
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 &= \overline{(\Xi^{\rho} \partial_{\rho} A^{\mu})} \underline{\partial}_{\mu} B + \overline{A^{\mu}} \underline{\partial}_{\mu} (\overline{\Xi^{\rho}} \partial_{\rho} B) \\
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- Scalar transformation: $\overline{\Xi^{\rho} \partial_{\rho}} F = \overline{\Xi^{\rho}} \underline{\partial}_{\rho} \underline{d}F = \Xi \cdot \underline{d}F$
- Every form dual to $\overline{A^{\rho} \underline{\partial}_{\rho}}$ is of the form

$$B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} \underline{d}X^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k}$$
- More operators than $\overline{A^{\rho} \underline{\partial}_{\rho}}$ dual to B (\longrightarrow old problem for vectors)

Problem:

Old problem: $F \cdot A^\mu \cdot F'$ does not transform as vector

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New problem: Define $F \cdot \overline{A^\mu \partial_\mu} \cdot F'$, acting as:

$$(F \cdot \overline{A^\rho \partial_\rho} \cdot F') B \equiv F \cdot (\overline{A^\mu \partial_\mu} B) \cdot F' = F''$$

→ Space of objects dual to B contains not only $\overline{A^\mu \partial_\mu}$,
but also $F \cdot \overline{A^\mu \partial_\mu} \cdot F'$ → much bigger space!

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New problem: Define $F \cdot \overline{A^\mu \partial_\mu} \cdot F'$, acting as:

$$(F \cdot \overline{A^\rho \partial_\rho} \cdot F') B \equiv F \cdot (\overline{A^\mu \partial_\mu} B) \cdot F' = F''$$

→ Space of objects dual to B contains not only $\overline{A^\mu \partial_\mu}$,
but also $F \cdot \overline{A^\mu \partial_\mu} \cdot F'$ → much bigger space!

New problem 2: What objects are dual to $F \cdot \overline{A^\mu \partial_\mu} \cdot F'$?

$$(G \cdot B \cdot G')(F \cdot \overline{A^\mu \partial_\mu} \cdot F') = G \cdot F \cdot \overline{A^\mu \partial_\mu} B \cdot F' \cdot G'$$

Problem:

Old problem: $F \cdot A^\mu \cdot F'$ does not transform as vector

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New problem 3: What objects are dual to $(G \cdot B \cdot G')$?

New problem 4: ...

No clear solution at this point...

3. Conclusions

- $F = \sum_{k=0}^{\infty} \frac{1}{k!} f_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_k}$

$$A^\mu = \sum_{k=0}^{\infty} \frac{1}{k!} a_{\lambda_1 \dots \lambda_k}^\mu X^{\lambda_1} \dots X^{\lambda_k}$$

$$B = \sum_{k=1}^{\infty} \sum_{j=1}^k b^{(j)}_{\lambda_1 \dots \lambda_k} X^{\lambda_1} \dots X^{\lambda_{j-1}} \underline{d} X^{\lambda_j} X^{\lambda_{j+1}} \dots X^{\lambda_k}$$

$$C = \sum_{k=2}^{\infty} \sum_{\substack{i,j=1 \\ i \neq j}}^k c_{\mu_1 \dots \mu_k}^{(i,j)} X^{\mu_1} \dots \underline{d}_{(1)} X^{\mu_i} \dots \underline{d}_{(2)} X^{\mu_j} \dots X^{\mu_k}$$

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$$\delta A^\mu = \overline{\Xi^\rho \partial_\rho} A^\mu - \overline{A^\rho \partial_\rho} \Xi^\mu$$

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- de Rahm operator \underline{d}
 scalar product $(A \cdot B) \longrightarrow$ problems

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- Possible ways out:

- Give up Leibnitz rule?
- Enhance definition for vector & give up transformation rule?
- ...?